

ventricle, and the extravasation covering the fissure in the aorta, exactly marked, as they appeared to,

My Lord,

Your Lordship's

most obedient

and most humble servant,

Frank Nicholls.

LII. *Of the Irregularities in the planetary Motions, caused by the mutual Attraction of the Planets : In a Letter to Charles Morton, M. D. Secretary to the Royal Society, by Charles Walmesley, F. R. S. and Member of the Royal Academy of Sciences at Berlin, and of the Institute at Bologna.*

S I R,

Read Dec. 10,
1761.

Finding that the influence, which the primary planets have upon one another, to disturb mutually their motions, had been but little considered, I thought it a subject worthy of examination. The force of the sun, to disturb the moon's motion, flows from the general principle of *gravitation*, and has been fully ascertained, both by theory and observation; and it follows, from the

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same principle, that all the planets must act upon one another, proportionally to the quantities of matter contained in their bulk, and inverse ratio of the squares of their mutual distances; but as the quantity of matter contained in each of them, is but small when compared to that of the sun, so their action upon one another, is not so sensible as that of the sun upon the moon. Astronomers generally contented themselves with solely considering those inequalities of the planetary motions, that arise from the elliptical figure of their orbits; but as they have been enabled, of late years, by the perfection of their instruments, to make observations with much more accuracy than before, they have discovered other variations, which they have not, indeed, been able yet to settle, but which seem to be owing to no other cause, but the mutual attraction of those celestial bodies. In order, therefore, to assist the astronomers in distinguishing and fixing these variations, I shall endeavour to calculate their quantity, from the general law of gravitation, and reduce the result into tables, that may be consulted, whenever observations are made.

I offer to you, at present, the first part of such a theory, in which I have chiefly considered the effects produced by the actions of the earth and Venus upon each other. But the same propositions will likewise give, by proper substitutions, the effects of the other planets upon these two, or of these two upon the others. To obviate, in part, the difficulty of such intricate calculations, I have supposed the orbits of the earth and Venus to be originally circular, and to suffer no other alteration, but what is occasioned by their mutual attraction, and the attraction of the other planets.

planets. Where the forces of two planets are considerable, with respect to each other, as in the case of Jupiter and Saturn, it may be necessary, in such computations, to have regard to the excentricity of their orbits; and this may be reserved for a subject of future scrutiny. But the supposing the orbits of the earth and Venus to be circular, may, in the present case, be admitted, without difficulty, as the forces of these two planets are so small, and the excentricity of their orbits not considerable. On these grounds, therefore, I have computed the variations, which are the effects of the earth's action: first, the variation of Venus's distance from the sun; secondly, that of its place in the ecliptic; thirdly, the retrograde motion of Venus's nodes; and, fourthly, the variation of inclination of its orbit to the plane of the ecliptic.

The similar irregularities in the motion of the earth, occasioned by its gravitation to Venus, are here likewise computed: but it is to be observed, that the absolute quantity of these irregularities is not here given, it being impossible, at present, to do it; because the absolute force of Venus is not known to us. I have, therefore, stated that planet's force by supposition, and have, accordingly, computed the effects it must produce; with the view, that the astronomers may compare their observations with the motions so calculated, and, from thence, discover how much the real force differs from that which has been supposed. But the exact determination of the force of Venus must be obtained, by observations made on the sun's place, at such times, when the effect of the other planets is either null or known.

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The influence of Venus upon the earth being thus computed, that of the other planets upon the same, may likewise, hereafter, be considered: by which means, the different equations, that are to enter into the settling of the sun's apparent place, will be determined; the change of the position of the plane of the earth's orbit will also be known; and, consequently, the alteration that thence arises in the obliquity of the ecliptic, and in the longitude and latitude of the fixed stars. These matters of speculation are reserved for another occasion, in case what is here offered should deserve approbation.

I am glad to have it in my power to present you with this testimony of my gratitude for past favours, and of my respect for your distinguished merit; and it is with sincerity, I subscribe myself,

S I R,

Your very humble servant,

Bath,
Nov. 21, 1761.

Cha. Walmesley.

*De Inæqualitatibus quas in motibus Planetarum
generant ipsorum in se invicem actiones.*

QUoniam in theoriæ hujus decursu frequens erit usus fluentium quæ arcubus circuli, vel eorum sinibus, cosinibus, et sinibus versis, exprimuntur, idcirco lemma sequens, quod alibi olim tradidi, lubet hîc apponere.

LEMMA.

LEMMA.

Dato cosinu arcûs cujusvis, invenire cosinum et finum arcûs alterius qui sit ad priorem in ratione λ ad 1.

Detur c cosinus arcûs A ad radium 1, et sit arcus $B = \lambda A$, cujus cosinus dicatur t ; eritque, ut notum est, $\dot{A} = \frac{-c}{\sqrt{1-c^2}}$, atque $\dot{B} = \lambda \dot{A} = \frac{-t}{\sqrt{1-t^2}}$.

Ponatur $c = \frac{1+xx}{2x}$, et $t = \frac{1+yy}{2y}$, fietque $\dot{A} = \frac{\dot{x}}{x\sqrt{1-c^2}}$,

$B = \frac{\dot{y}}{y\sqrt{1-t^2}}$: sed est $\dot{A} : \dot{B} :: 1 : \lambda$, adeoque $\frac{\lambda \dot{x}}{x} = \frac{\dot{y}}{y}$,

unde $\log. x^\lambda = \log. y$, et $x^\lambda = y$. Verùm æquationes

$c = \frac{1+xx}{2x}$ et $t = \frac{1+yy}{2y}$ dant $x = c + \sqrt{cc-1}$,

$x = c - \sqrt{cc-1}$, et $y = t + \sqrt{tt-1}$, $y =$

$t - \sqrt{tt-1}$; unde est $x^\lambda = t + \sqrt{tt-1} =$

$\frac{c + \sqrt{cc-1}}{c - \sqrt{cc-1}}^\lambda$, atque inde $2t = c + \sqrt{cc-1}^\lambda$

$+ c - \sqrt{cc-1}^\lambda$. Fiat igitur $c + \sqrt{cc-1} = l$,

et $c - \sqrt{cc-1} = m$, eritque $lm = 1$, et $c = \cos.$

$A = \frac{l+m}{2}$, et $\sin. A = \frac{l-m}{2} \sqrt{-1}$; atque inde

$t = \cos. B = \frac{l^\lambda + m^\lambda}{2}$, et $\sin. B = \frac{l^\lambda - m^\lambda}{2} \sqrt{-1}$.

Itaque in circulo, cujus radius est 1, si duorum arcuum vel angulorum A et B alteruter B sit ad alterum A ut numerus quilibet λ ad 1, et ponatur

$\cos. A = \frac{l+m}{2}$, existente $lm = 1$, erit $\sin. A$

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$$= \frac{l-m}{2} \sqrt{-1}, \text{ atque } \cos. B = \cos. \lambda A = \frac{l^{\lambda} + m^{\lambda}}{2},$$

et $\sin. B = \sin. \lambda A = \frac{l^{\lambda} - m^{\lambda}}{2} \sqrt{-1}. \quad Q. E. I.$

COROLL. I.

Hinc habetur $\cos. A \times \cos. B = \frac{l+m}{2} \times \frac{l^{\lambda} + m^{\lambda}}{2} =$
 $\frac{l^{\lambda+1} + m^{\lambda+1}}{4} + \frac{l^{\lambda-1} + m^{\lambda-1}}{4};$ sed, quemadmodum per
hoc lemma est $\frac{l^{\lambda} + m^{\lambda}}{2} = \cos. \lambda A$, erit $\frac{l^{\lambda+1} + m^{\lambda+1}}{2} =$
 $\cos. \overline{\lambda + 1} \times A = \cos. A + B$, atque $\frac{l^{\lambda-1} + m^{\lambda-1}}{2} =$
 $\cos. \overline{\lambda - 1} \times A = \cos. \overline{B - A}$, adeoque $\cos. A \times \cos. B$
 $= \frac{1}{2} \cos. \overline{A + B} + \frac{1}{2} \cos. \overline{B - A}.$

Atque hoc calculi methodo facilè eruuntur sequentes formulæ pro duobus angulis A et B, advertendo esse $\cos. \overline{B - A} = \cos. \overline{A - B}$, $\sin. \overline{B - A} = -\sin. \overline{A - B}$, et $\cos. 0 = 1$.

1°. $\cos. A \times \cos. B = \frac{1}{2} \cos. \overline{A + B} + \frac{1}{2} \cos. \overline{A - B}.$

2°. $\sin. A \times \sin. B = -\frac{1}{2} \cos. \overline{A + B} + \frac{1}{2} \cos. \overline{A - B}.$

3°. $\sin. A \times \cos. B = \frac{1}{2} \sin. \overline{A + B} + \frac{1}{2} \sin. \overline{A - B}.$

Atque ex illis hæ sequentes eliciuntur,

4°. $\cos. \overline{A + B} = \cos. A \times \cos. B - \sin. A \times \sin. B.$

5°. $\cos. A - B = \sin. A \times \sin. B + \cos. A \times \cos. B.$

6°. $\sin. A + B = \sin. A \times \cos. B + \cos. A \times \sin. B.$

7°. $\sin. A - B = \sin. A \times \cos. B - \cos. A \times \sin. B.$

Tùm ex his valores tangentium haud ægrè derivantur,

Quippe

Quippe cum fit generatim pro quovis angulo A,

$$\text{tang. } A = \frac{\text{fin. } A}{\text{cof. } A}, \text{ erit } \text{tang. } \overline{A + B} = \frac{\text{fin. } \overline{A + B}}{\text{cof. } \overline{A + B}} =$$

$$\frac{\text{fin. } A \times \text{cof. } B + \text{cof. } A \times \text{fin. } B}{\text{cof. } A \times \text{cof. } B - \text{fin. } A \times \text{fin. } B} = \frac{\text{fin. } A \times \text{cof. } B + \text{cof. } A \times \text{fin. } B}{\text{cof. } A \times \text{fin. } B}.$$

$$\times \frac{\text{cof. } A \times \text{fin. } B}{\text{cof. } A \times \text{cof. } B - \text{fin. } A \times \text{fin. } B} = \frac{\text{tang. } A}{\text{tang. } B} + 1$$

$$\times \frac{1}{\frac{1}{\text{tang. } B} - \text{tang. } A} = \frac{\text{tang. } A + \text{tang. } B}{1 - \text{tang. } A \times \text{tang. } B}. \text{ Simili}$$

$$\text{calculo prodit } \text{tang. } A - B = \frac{\text{tang. } A - \text{tang. } B}{1 + \text{tang. } A \times \text{tang. } B}.$$

Unde statui possunt,

$$1^{\circ}. \text{Tang. } \overline{A + B} = \frac{\text{tang. } A + \text{tang. } B}{1 - \text{tang. } A \times \text{tang. } B}.$$

$$2^{\circ}. \text{Tang. } \overline{A - B} = \frac{\text{tang. } A - \text{tang. } B}{1 + \text{tang. } A \times \text{tang. } B}.$$

$$3^{\circ}. \text{Tang. } A \times \text{tang. } B = \frac{\text{tang. } \overline{A + B} - \text{tang. } A - \text{tang. } B}{\text{tang. } A + B},$$

$$\text{vel } \text{tang. } A \times \text{tang. } B = \frac{\text{tang. } A - \text{tang. } B - \text{tang. } \overline{A - B}}{\text{tang. } A - B}.$$

COROLL. II.

Erat in lemmate $\dot{A} = \frac{x}{x \sqrt{-1}}$, unde est $A \sqrt{-1} = \log. x$.

Denotet igitur E numerum cujus logarithmus hyperbolicus est 1, eritque $E^{A \sqrt{-1}} = x$, et cum sit $x = c + \sqrt{cc - 1}$, inde obtinetur $c = \text{cof. } A = \frac{E^{A \sqrt{-1}} + E^{-A \sqrt{-1}}}{2}$, atque $\text{fin. } A = \frac{E^{A \sqrt{-1}} - E^{-A \sqrt{-1}}}{2 \sqrt{-1}}$.

Sunt qui his finuum et cosinuum valoribus potiùs utuntur; verùm ii valores, quos exhibet corollarium præcedens, simpliciores sunt et calculo plerumque aptiores.

COROLL. III.

Quoniam est $2 \times \text{cos. } A = l + m$, erit

$$2^\lambda \times \overline{\text{cos. } A}^\lambda = \begin{cases} l^\lambda + \lambda l^{\lambda-1} m + \lambda \times \frac{\lambda-1}{2} l^{\lambda-2} m^2 + \lambda \\ \times \frac{\lambda-1}{2} \times \frac{\lambda-2}{3} l^{\lambda-3} m^3 +, \&c. \\ m^\lambda + \lambda m^{\lambda-1} l + \lambda \times \frac{\lambda-1}{2} m^{\lambda-2} l^2 + \lambda \\ \times \frac{\lambda-1}{2} \times \frac{\lambda-2}{3} m^{\lambda-3} l^3 +, \&c. \end{cases}$$

assumendo scilicet primos et ultimos terminos homologos seriei exprimentis quantitatem $\overline{l+m}^\lambda$: unde, propter $lm = 1$, provenit

$$2^{\lambda-1} \times \overline{\text{cos. } A}^\lambda = \frac{l^\lambda + m^\lambda}{2} + \lambda \times \frac{l^{\lambda-2} + m^{\lambda-2}}{2} + \lambda \times \frac{\lambda-1}{2} \\ \times \frac{l^{\lambda-4} + m^{\lambda-4}}{2} + \lambda \times \frac{\lambda-1}{2} \times \frac{\lambda-2}{3} \times \frac{l^{\lambda-6} + m^{\lambda-6}}{2} +, \&c.$$

atque adeò per lemma

$$\overline{\text{cos. } A}^\lambda = \frac{1}{2^{\lambda-1}} \text{in cos. } \lambda A + \lambda \text{ cos. } \overline{\lambda-2} \times A + \lambda \\ \times \frac{\lambda-1}{2} \text{cos. } \overline{\lambda-4} \times A + \lambda \times \frac{\lambda-1}{2} \times \frac{\lambda-2}{3} \text{cos. } \overline{\lambda-6} \\ \times A +, \&c.$$

Ubi λ est numerus impar, terminus ultimus seriei erit ille in quo numerus λ , vel $\lambda - 2$, vel $\lambda - 4$, &c. qui multiplicat angulum A , evadit æqualis 1. Ubi verò λ est numerus par, terminus ultimus seriei erit ille in quo numerus prædictus evadit æqualis 0,

quo in casu semiffis tantum ultimi termini sumenda est; cum enim series hæc colligatur ex numero pari terminorum homologorum, quæ tamen, ubi λ est numerus par, constare debet ex terminorum numero impari, idè duplum exhibet terminum ultimum.

Simili modo cum fit $2 \times \sin. A = \sqrt{1-m} \times \sqrt{1-I}$, erit

$$2^{\lambda} \times \overline{\sin. A}^{\lambda} = \sqrt{1-I} \times \left\{ \begin{array}{l} I^{\lambda} - \lambda I^{\lambda-1} m + \lambda \times \frac{\lambda-1}{2} I^{\lambda-2} m^2 \\ - \lambda \times \frac{\lambda-1}{2} \times \frac{\lambda-2}{3} I^{\lambda-3} m^3 +, \\ \&c. \\ + m^{\lambda} + \lambda m^{\lambda-1} I + \lambda \times \frac{\lambda-1}{2} m^{\lambda-2} I^2 \\ + \lambda \times \frac{\lambda-1}{2} \times \frac{\lambda-2}{3} m^{\lambda-3} I^3 +, \\ \&c. \end{array} \right.$$

Terminis inferioribus hujus seriei præfiguntur alternatim signa $+$ $-$ ubi λ est numerus par, et signa $-$ $+$ ubi λ est numerus impar, adeoque in priore casu est

$$2^{\lambda-1} \times \overline{\sin. A}^{\lambda} = \sqrt{1-I}^{\lambda} \text{ in } \frac{I^{\lambda} + m^{\lambda}}{2} - \lambda \times \frac{I^{\lambda-2} + m^{\lambda-2}}{2} + \lambda \times \frac{\lambda-1}{2} \times \frac{I^{\lambda-4} + m^{\lambda-4}}{2} - \lambda \times \frac{\lambda-1}{2} \times \frac{\lambda-2}{3} \times \frac{I^{\lambda-6} + m^{\lambda-6}}{2} +, \&c.$$

et in casu posteriori

$$2^{\lambda-1} \times \overline{\sin. A}^{\lambda} = \sqrt{1-I}^{\lambda} \text{ in } \frac{I^{\lambda} - m^{\lambda}}{2} - \lambda \times \frac{I^{\lambda-2} - m^{\lambda-2}}{2} + \lambda \times \frac{\lambda-1}{2} \times \frac{I^{\lambda-4} - m^{\lambda-4}}{2} - \lambda \times \frac{\lambda-1}{2} \times \frac{\lambda-2}{3} \times \frac{I^{\lambda-6} - m^{\lambda-6}}{2} +, \&c.$$

Adeoque si λ sit numerus par, erit

$$\overline{\sin. A}^{\lambda} = \frac{I}{2^{\lambda-1}} \text{ in } \pm \cos. \lambda A \mp \lambda \cos. \frac{\lambda-2}{2} \times A \pm \lambda \times \frac{\lambda-1}{2} \cos. \frac{\lambda-4}{2} \times A \mp \lambda \times \frac{\lambda-1}{2} \times \frac{\lambda-2}{3} \cos. \frac{\lambda-6}{2} \times A \pm, \&c.$$

Signa hęc alternatim mutantur, et superiora sunt adhibenda, ubi λ exprimit unum ex numeris 4, 8, 12, 16, &c. quia tunc est $\sqrt{-1}^\lambda = 1$; inferiora autem adhibenda, ubi λ exprimit unum ex numeris 2, 6, 10, 14, &c. quia tunc est $\sqrt{-1}^\lambda = -1$.

Si λ fit numerus impar, cū per lemma fit $\frac{l^\lambda - m^\lambda}{2} \sqrt{-1} = \text{fin. } \lambda A$, et $\frac{l^{\lambda-2} - m^{\lambda-2}}{2} \sqrt{-1} = \text{fin. } \lambda - 2 \times A$, &c. habetur

$$\begin{aligned} \overline{\text{fin. } A}^\lambda &= \frac{1}{2^{\lambda-1}} \text{in } \pm \text{fin. } \lambda A \mp \lambda \times \text{fin. } \lambda - 2 \times A \pm \lambda \\ &\times \frac{\lambda-1}{2} \text{fin. } \lambda - 4 \times A \mp \lambda \times \frac{\lambda-1}{2} \times \frac{\lambda-2}{3} \text{fin. } \lambda - 6 \\ &\times A \pm, \text{ \&c.} \end{aligned}$$

ubi signa superiora sunt usurpanda, cū λ exprimit unum ex numeris 1, 5, 9, 13, &c. quia tunc est $\sqrt{-1}^\lambda = \sqrt{-1}$; et signa inferiora, cū λ fuerit unus ex numeris 3, 7, 11, 15, &c. quia tunc est $\sqrt{-1}^\lambda = -\sqrt{-1}$.

Notandum autem, seriei ultimum terminum esse illum in quo numerus λ , vel $\lambda - 2$, vel $\lambda - 4$, &c. est æqualis 1 ubi λ est numerus impar; atque terminum ultimum esse illum in quo prædictus numerus est æqualis 0 ubi λ est numerus par, quo in casu semissis tantū ultimi termini assumenda est ob rationem superiùs datam.

Ex his finuum et cosinuum expressionibus alia huiusmodi theoremata deducere liceret, sed quæ hęc traduntur ad præsens institutum sufficiunt.

COROLL. IV.

Notum est fluentem fluxionis \dot{A} $\cos. A$ esse $\sin. A$,
 atque fluentem fluxionis \dot{A} $\sin. A$ esse $\sin. \text{verf. } A$.
 Pariter si fumatur arcus λA qui fit ad arcum A ut
 numerus quilibet λ ad 1, cùm sit $\lambda \dot{A} \cos. \lambda A$ æqua-
 lis fluxioni finûs arcûs λA , erit flu. $\dot{A} \cos. \lambda A =$
 $\frac{\sin. \lambda A}{\lambda}$, et flu. $\dot{A} \sin. \lambda A = \frac{\sin. \text{verf. } \lambda A}{\lambda}$. Itemque, si
 ad arcum λA adjungatur arcus datus d , cùm fluxio
 arcûs $\lambda A + d$ sit æqualis $\lambda \dot{A}$, erit flu. $\dot{A} \cos. \lambda A + d$
 $= \frac{\sin. \overline{\lambda A + d}}{\lambda}$, et flu. $\dot{A} \sin. \overline{\lambda A + d} = \frac{\sin. \text{verf. } \overline{\lambda A + d}}{\lambda}$.

Sumantur jam duo anguli, vel duo arcus λA et μA ,
 qui sint ad angulum, vel arcum A respectivè, ut λ et
 μ ad 1, atque per Coroll. II. habetur $\cos. \lambda A \cos. \mu A$
 $= \frac{1}{2} \cos. \overline{\lambda + \mu} \times A + \frac{1}{2} \cos. \overline{\lambda - \mu} \times A$; unde
 erit fluens fluxionis $\dot{A} \cos. \lambda A \times \cos. \mu A$ æqualis
 $\frac{\sin. \overline{\lambda + \mu} \times A}{2 \times \lambda + \mu} + \frac{\sin. \overline{\lambda - \mu} \times A}{2 \times \lambda - \mu}$.

Atque hoc methodo prodeunt sequentes formulæ

$$1^{\circ}. \text{ Flu. } \dot{A} \cos. \lambda A \times \cos. \mu A = \frac{\sin. \overline{\lambda + \mu} \times A}{2 \times \lambda + \mu} + \frac{\sin. \overline{\lambda - \mu} \times A}{2 \times \lambda - \mu}.$$

$$2^{\circ}. \text{ Flu. } \dot{A} \sin. \lambda A \times \sin. \mu A = - \frac{\sin. \overline{\lambda + \mu} \times A}{2 \times \lambda + \mu} + \frac{\sin. \overline{\lambda - \mu} \times A}{2 \times \lambda - \mu}.$$

3^o. Flu.

$$3^{\circ}. \text{Flu. } \dot{A} \text{ fin. } \lambda A \times \text{cof. } \mu A = \frac{\text{fin. verf. } \overline{\lambda + \mu} \times A}{2 \times \lambda + \mu} \\ + \frac{\text{fin. verf. } \overline{\lambda - \mu} \times A}{2 \times \lambda - \mu}.$$

Advertendum autem est, ubi $\lambda = \mu$, tunc esse
 $\text{cof. } \lambda A \times \text{cof. } \mu A = \frac{1}{2} \text{cof. } 2\lambda A + \frac{1}{2}$, $\text{fin. } \lambda A$
 $\times \text{fin. } \mu A = -\frac{1}{2} \text{cof. } 2\lambda A + \frac{1}{2}$, $\text{fin. } \lambda A \times \text{cof. } \mu A$
 $= \frac{1}{2} \text{fin. } 2\lambda A$; adeoque in hoc casu formulæ præ-
cedentes evadunt

$$1^{\circ}. \text{Flu. } \dot{A} \times \overline{\text{cof. } \lambda A}^2 = \frac{\text{fin. } 2\lambda A}{4\lambda} + \frac{A}{2}.$$

$$2^{\circ}. \text{Flu. } \dot{A} \times \overline{\text{fin. } \lambda A}^2 = -\frac{\text{fin. } 2\lambda A}{4\lambda} + \frac{A}{2}.$$

$$3^{\circ}. \text{Flu. } \dot{A} \times \text{fin. } \lambda A \times \text{cof. } \lambda A = \frac{\text{fin. verf. } 2\lambda A}{4\lambda}.$$

Si angulo λA addatur angulus datus d , erit cof.
 $\overline{\lambda A + d} \times \text{cof. } \mu A = \frac{1}{2} \text{cof. } \overline{\lambda + \mu} \times A + d + \frac{1}{2}$
 $\text{cof. } \overline{\lambda - \mu} \times A + d$, atque inde

$$1^{\circ}. \text{Flu. } \dot{A} \text{ cof. } \overline{\lambda A + d} \times \text{cof. } \mu A = \frac{\text{fin. } \overline{\lambda + \mu} \times A + d}{2 \times \lambda + \mu} \\ + \frac{\text{fin. } \overline{\lambda - \mu} \times A + d}{2 \times \lambda - \mu}.$$

$$2^{\circ}. \text{Flu. } \dot{A} \text{ fin. } \overline{\lambda A + d} \times \text{fin. } \mu A = -\frac{\text{fin. } \overline{\lambda + \mu} \times A + d}{2 \times \lambda + \mu} \\ + \frac{\text{fin. } \overline{\lambda - \mu} \times A + d}{2 \times \lambda - \mu}.$$

$$3^{\circ}. \text{Flu. } \dot{A} \text{ fin. } \overline{\lambda A + d} \times \text{cof. } \mu A = \frac{\text{fin. verf. } \overline{\lambda + \mu} \times A + d}{2 \times \lambda + \mu} \\ + \frac{\text{fin. verf. } \overline{\lambda - \mu} \times A + d}{2 \times \lambda - \mu}.$$

4^o. Flu.

$$4^{\circ}. \text{Flu. } \dot{A} \text{ cof. } \overline{\lambda A + d} \times \text{fin. } \mu A = \frac{\text{fin. verf. } \overline{\lambda + \mu \times A + d}}{2 \times \lambda + \mu} \\ - \frac{\text{fin. verf. } \overline{\lambda - \mu \times A + d}}{2 \times \lambda - \mu}.$$

Si fuerit $\lambda = \mu$, erit $\text{cof. } \overline{\lambda A + d} \times \text{cof. } \lambda A = \frac{1}{2} \text{cof. } 2\lambda A + d + \frac{1}{2} \text{cof. } d$, &c. adeoque formulæ præcedentes in has abeunt,

$$1^{\circ}. \text{Flu. } \dot{A} \text{ cof. } \overline{\lambda A + d} \times \text{cof. } \lambda A = \frac{\text{fin. } \overline{2\lambda A + d}}{4\lambda} \\ + \frac{\text{cof. } d}{2} A.$$

$$2^{\circ}. \text{Flu. } \dot{A} \text{ fin. } \overline{\lambda A + d} \times \text{fin. } \lambda A = - \frac{\text{fin. } \overline{2\lambda A + d}}{4\lambda} \\ + \frac{\text{cof. } d}{2} A.$$

$$3^{\circ}. \text{Flu. } \dot{A} \text{ fin. } \overline{\lambda A + d} \times \text{cof. } \lambda A = \frac{\text{fin. verf. } \overline{2\lambda A + d}}{4\lambda} \\ + \frac{\text{fin. } d}{2} A.$$

$$4^{\circ}. \text{Flu. } \dot{A} \text{ cof. } \overline{\lambda A + d} \times \text{fin. } \lambda A = \frac{\text{fin. verf. } \overline{2\lambda A + d}}{4\lambda} \\ - \frac{\text{fin. } d}{2} A.$$

PROPOSITIO I. PROBLEMA.

In systemate duorum planetarum circa Solem in orbibus penè circularibus revolventium, requiratur vis planetæ exterioris ad perturbandum motum interioris.

Revolvantur planetæ duo P et Q (Fig. I.) in eodem plano circa Solem in S, et jungantur SP, SQ, PQ.

Orbis

planetarum referantur ad Solem spectatum tanquàm immotum, vis $\frac{\phi k}{z^3}$ pars ea $\frac{\phi}{k^2}$, qua simul urgentur Sol et planeta Q versus P secundum lineas parallelas, non mutat corporum S et Q situm ad se invicem, ideoque differentia virium sola perturbationem inducit.

Quare differentia illa, nimirum $\frac{\phi k}{z^3} - \frac{\phi}{k^2}$, exponatur per lineam QT parallelam rectæ PS, et in SQ demisso perpendiculo TR, vis QT resolvetur in vires TR, QR, eritque vis QT ad vim TR ut radius 1 ad finum anguli QSP, adeoque vis TR = $\frac{\phi k}{z^3} - \frac{\phi}{k^2}$

× sin. QSP, et vis QR = $\frac{\phi k}{z^3} - \frac{\phi}{k^2}$ × cos. QSP. Ex vi autem QR tollatur vis OI utpotè in contrarium agens, et manebit vis $\frac{\phi k}{z^3} - \frac{\phi}{k^2}$ × cos. QSP = $\frac{\phi}{z^3}$.

Vires igitur, quibus planeta P perturbat motum planetæ Q quatenus in eodem plano moventur, sunt

1°. Vis TR ad radium QS perpendicularis, qua augetur vel minuitur area tempore dato descripta, estque æqualis $\frac{\phi k}{z^3} - \frac{\phi}{k^2}$ × sin. QSP.

2°. Vis $\frac{\phi}{z^3} \times k \cos. QSP - 1 - \frac{\phi}{k^2} \cos. QSP$, qua retrahitur planeta Q à Sole in directione radii SQ.

Ut autem harumce virium expressiones formam induant calculo accommodam, ope trianguli PSQ habebitur $PQ^2 = z^2 = kk + xx - 2kx \times \cos. QSP$, five, posita $x = 1$ ob rationem dictam, $z^2 = kk$

+ 1 — 2k × cof. QSP. Affumatur jam angulus θ qui semper sit ad angulum QSP in ratione n ad 1, eritque $QSP = \frac{1}{n}s$, et posito $kk + 1 = tt$, et $\frac{2k}{t^2} =$

b , erit $x^2 = t^2 \times 1 - b \text{ cof. } \frac{1}{n}s$, hincque $\frac{1}{x^2} = \frac{1}{t^2}$

$\times 1 - b \text{ cof. } \frac{1}{n}s \Big]^{-\frac{1}{2}}$. Si b fuerit unitati ferè æqualis,

et evolvatur quantitas $1 - b \text{ cof. } \frac{1}{n}s \Big]^{-\frac{1}{2}}$ in seriem modo

solito, series illa parùm convergit, estque ad operationes analyticas minùs commoda. Series igitur alia investiganda est, et quia ex lemmate patet hujusmodi quantitatem $\overline{\text{cof. A}}^n$ exprimi posse aggregato terminorum, quorum singuli ducuntur in cosinus angulorum qui sunt anguli A multiplices, generatim sup-

ponemus $1 - b \text{ cof. } \frac{1}{n}s \Big]^n = R + S \text{ cof. } \frac{1}{n}s + T \text{ cof. } \frac{2}{n}s$

$+ V \text{ cof. } \frac{3}{n}s + W \text{ cof. } \frac{4}{n}s + \&c.$

Atque ut inveniantur valores coefficientium R, S,

T, &c. fumatur utrinque fluxio, nempe $\frac{mb}{n}s \times \text{fin. } \frac{1}{n}s$

$\times 1 - b \text{ cof. } \frac{1}{n}s \Big]^{n-1} = -S \times \frac{1}{n}s \times \text{fin. } \frac{1}{n}s - T \times \frac{2}{n}s$

$\times \text{fin. } \frac{2}{n}s - V \times \frac{3}{n}s \times \text{fin. } \frac{3}{n}s - W \times \frac{4}{n}s \times \text{fin. } \frac{4}{n}s -$

&c. atque ducatur æquatio hæc in $1 - b \text{ cof. } \frac{1}{n}s$, et

substituto pro $1 - b \text{ cof. } \frac{1}{n}s \Big]^n$ ipsius valore $R + S \text{ cof. } \frac{1}{n}s$

$+ T \text{ cof. } \frac{2}{n}s + \&c.$ fiet $mb \times \text{fin. } \frac{1}{n}s$

×

$$*R + S \operatorname{cof.} \frac{1}{n} s + T \operatorname{cof.} \frac{2}{n} s + V \operatorname{cof.} \frac{3}{n} s +, \&c. = 1 - b \operatorname{cof.} \frac{1}{n} s$$

$$* - S \times \sin. \frac{1}{n} s + 2 T \times \sin. \frac{2}{n} s + 3 V \times \sin. \frac{3}{n} s + W \times \sin. \frac{4}{n} s +, \&c.$$

et factâ multiplicatione, cùm fit (per Coroll. I. Lem.)

$$\sin. \frac{1}{n} s \times \operatorname{cof.} \frac{r}{n} s = \frac{1}{2} \sin. \frac{r+1}{n} s - \frac{1}{2} \sin. \frac{r-1}{n} s, \text{ ac}$$

$$\sin. \frac{r}{n} s \times \operatorname{cof.} \frac{1}{n} s = \frac{1}{2} \sin. \frac{r+1}{n} s + \frac{1}{2} \sin. \frac{r-1}{n} s, \text{ emerget}$$

$$\left. \begin{array}{l} + 2mbR \\ + 2S \\ - 2bT \\ - mbT \end{array} \right\} \times \sin. \frac{1}{n} s \left\{ \begin{array}{l} + mbS \\ - bS \\ + 4T \\ - 3bV \\ - mbV \end{array} \right\} \times \sin. \frac{2}{n} s \left\{ \begin{array}{l} + mbT \\ - 2bT \\ + 6V \\ - 4bW \\ - mbW \end{array} \right\} \times \sin. \frac{3}{n} s, \&c. = 0$$

Deinde nihilo æquando singulos terminos, prodeunt

$$T = \frac{2S + 2mbR}{m + 2 \times b}, \quad V = \frac{4T + \overline{m-1} \times bS}{m + 3 \times b}, \quad W =$$

$$\frac{6V + \overline{m-2} \times bT}{m + 4 \times b}, \&c. \text{ quorum valorum progressus satis manifestus est.}$$

Datis igitur primis duobus coefficientibus R et S, dabuntur et reliqui: R et S autem sic inveniuntur.

$$\begin{aligned} \text{Est } 1 - b \operatorname{cof.} \frac{1}{n} s &= 1 - mb \operatorname{cof.} \frac{1}{n} s + m \\ &\times \frac{m-1}{2} b^2 \operatorname{cof.} \frac{1}{n} s^2 - m \times \frac{m-1}{2} \times \frac{m-2}{3} b^3 \operatorname{cof.} \frac{1}{n} s^3, \\ \&c. &= R + S \operatorname{cof.} \frac{1}{n} s + T \operatorname{cof.} \frac{2}{n} s + V \operatorname{cof.} \frac{3}{n} s +, \\ \&c. &\text{Evolvantur termini } \operatorname{cof.} \frac{1}{n} s^2, \operatorname{cof.} \frac{1}{n} s^4, \operatorname{cof.} \frac{1}{n} s^6, \\ \&c. &\text{per methodum traditam in Coroll. III. Lem.} \\ \text{ac, collectis simul omnibus terminis qui nullo cosinu} &\text{afficiuntur, prodibit} \end{aligned}$$

$$R = 1 + \frac{m}{2} \times \frac{m-1}{2} b^2 + \frac{m}{2} \times \frac{m-1}{2} \times \frac{m-2}{4} \times \frac{m-3}{4} b^4 \\ + \frac{m}{2} \times \frac{m-1}{2} \times \frac{m-2}{4} \times \frac{m-3}{4} \times \frac{m-4}{6} \times \frac{m-5}{6} b^6 +, \&c.$$

cujus seriei progressio satis patet; atque adeò, cum sit in hoc nostro problemate $m = -\frac{3}{2}$, erit

$$R = 1 + \frac{3 \times 5}{4 \times 4} b^2 + \frac{3 \times 5}{4 \times 4} \times \frac{7 \times 9}{8 \times 8} b^4 + \frac{3 \times 5}{4 \times 4} \times \frac{7 \times 9}{8 \times 8} \\ \times \frac{11 \times 13}{12 \times 12} b^6 + \frac{3 \times 5}{4 \times 4} \times \frac{7 \times 9}{8 \times 8} \times \frac{11 \times 13}{12 \times 12} \times \frac{15 \times 17}{16 \times 16} b^8 +, \&c.$$

Insipienti indolem hujus seriei patebit terminum quemlibet æquari termino antecedenti ducto in

$$\frac{r+1 \times r-1}{r^2} b^2, \text{ five } \frac{r^2-1}{r^2} b^2, \text{ } r \text{ existente æquali nu-}$$

mero quadruplicato terminorum præcedentium: sic terminus sextus, quia habetur in hoc casu $r = 5 \times 4$

$$= 20, \text{ æqualis est termino quinto } \frac{3 \times 5}{4 \times 4} \dots \frac{15 \times 17}{16 \times 16} b^8$$

$$\text{ducto in } \frac{19 \times 21}{20 \times 20} b^2.$$

Termino igitur quovis hujus seriei dicto B, terminus subsequens erit $Bb^2 \times \frac{r^2-1}{r^2}$: et manente de-

inceps eodem, quem in hoc termino habet, numeri r valore, termini subsequentes erunt, $Bb^4 \times \frac{r^2-1}{r^2}$

$$\times \frac{(r+4)^2-1}{(r+4)^2}, Bb^6 \times \frac{r^2-1}{r^2} \times \frac{(r+4)^2-1}{(r+4)^2} \times \frac{(r+8)^2-1}{(r+8)^2},$$

$$Bb^8 \times \frac{r^2-1}{r^2} \dots \frac{(r+12)^2-1}{(r+12)^2}, \&c. \text{ Sed est } \frac{r^2-1}{r^2} =$$

$$1 - \frac{1}{r^2}, \frac{(r+4)^2-1}{(r+4)^2} = 1 - \frac{1}{(r+4)^2}, \&c. \text{ et si fuerit } r$$

numerus

numerus aliquantùm magnus, erit $\frac{r^2 - 1}{r^2} \times \frac{\overline{r+4}^2 - 1}{(r+4)^2}$

$$= 1 - \frac{1}{r^2} - \frac{1}{(r+4)^2}, \text{ et } \frac{r^2 - 1}{r^2} \times \frac{\overline{r+4}^2 - 1}{(r+4)^2} \times \frac{\overline{r+8}^2 - 1}{(r+8)^2}$$

$$= 1 - \frac{1}{r^2} - \frac{1}{(r+4)^2} - \frac{1}{(r+8)^2}, \text{ atque ita porro, rejiciendo}$$

fractiones hujus generis $\frac{1}{r^2 \times (r+4)^2}$ et alias his minores.

Unde termini omnes prædicti, incipiendo à termino B, erunt

$$\begin{aligned} B + Bb^2 + Bb^4 + Bb^6 + Bb^8 +, \&c. &= B \times \frac{1}{1-b^2} \\ - \frac{Bb^2}{r^2} - \frac{Bb^4}{r^2} - \frac{Bb^6}{r^2} - \frac{Bb^8}{r^2} -, \&c. &= - \frac{B}{r^2} \times \frac{b^2}{1-b^2} \\ - \frac{Bb^4}{(r+4)^2} - \frac{Bb^6}{(r+4)^2} - \frac{Bb^8}{(r+4)^2} -, \&c. &= - \frac{B}{(r+4)^2} \times \frac{b^4}{1-b^2} \\ - \frac{Bb^6}{(r+8)^2} - \frac{Bb^8}{(r+8)^2} -, \&c. &= - \frac{B}{(r+8)^2} \times \frac{b^6}{1-b^2} \\ - \frac{Bb^8}{(r+12)^2} -, \&c. &= - \frac{B}{(r+12)^2} \times \frac{b^8}{1-b^2} \\ \&c. &\&c. \end{aligned}$$

ac proinde tandem fit

$$\begin{aligned} R &= 1 + \frac{3 \times 5}{4 \times 4} b^2 + \frac{3 \times 5}{4 \times 4} \times \frac{7 \times 9}{8 \times 8} b^4 + \frac{3 \times 5}{4 \times 4} \times \frac{7 \times 9}{8 \times 8} \\ &\times \frac{11 \times 13}{12 \times 12} b^6 \times \frac{3 \times 5}{4 \times 4} \dots \frac{15 \times 17}{16 \times 16} b^8 +, \&c. + \frac{B}{1-b^2} \\ &\times 1 - \frac{b^2}{r^2} - \frac{b^4}{(r+4)^2} - \frac{b^6}{(r+8)^2} - \frac{b^8}{(r+12)^2} - \frac{b^{10}}{(r+16)^2} -, \&c. \end{aligned}$$

Unde si, computatis, exempli gratiâ, decem terminis, undecimus designetur per B, erit $r = 10 \times 4 = 40$, et summa illorum decem terminorum addita summæ

summæ seriei $\frac{B}{1-b^2} \times 1 - \frac{b^2}{1-b^4} - \frac{b^4}{1-b^8} - \dots$, &c. dabit
valorem ipsius R.

Simili modo si in æquatione prædictâ $1 - mb \cos. \frac{s}{n}$
 $+ m \times \frac{m-1}{2} b^2 \times \cos. \frac{1}{n} s^2 + m \times \frac{m-1}{2} \times \frac{m-2}{3} b^3$
 $\times \cos. \frac{1}{n} s^3 + \dots = R + S \cos. \frac{s}{n} + T \cos. \frac{2}{n} s$
 $+ V \cos. \frac{3}{n} s + \dots$ evolvantur quantitates $\cos. \frac{1}{n} s$,
 $\cos. \frac{1}{n} s^2$, $\cos. \frac{1}{n} s^3$, &c. in suos valores, prout in
 Coroll. III. Lem. edoctum est, et colligantur omnes
 termini qui ducuntur in $\cos. \frac{1}{n} s$ exurget

$$S = -mb - m \times \frac{m-1}{2} \times \frac{m-2}{4} b^3 - m \times \frac{m-1}{2} \\ \times \frac{m-2}{4} \times \frac{m-3}{4} \times \frac{m-4}{6} b^5 - m \times \frac{m-1}{2} \times \frac{m-2}{4} \\ \times \frac{m-3}{4} \times \frac{m-4}{6} \times \frac{m-5}{6} \times \frac{m-6}{8} b^7 - \dots$$

five, posito $m = -\frac{3}{2}$,

$$S = \frac{3}{2} b + \frac{3}{2} \times \frac{5 \times 7}{4 \times 8} b^3 + \frac{3}{2} \times \frac{5 \times 7}{4 \times 8} \times \frac{9 \times 11}{8 \times 12} b^5 + \frac{3}{2} \\ \times \frac{5 \times 7}{4 \times 8} \times \frac{9 \times 11}{8 \times 12} \times \frac{13 \times 15}{12 \times 16} b^7 + \frac{3}{2} \dots \frac{13 \times 15}{12 \times 16} \\ \times \frac{17 \times 19}{16 \times 20} b^9 + \dots$$

Patet autem terminum quemlibet hujus seriei æquari
 termino antecedenti ducto in $\frac{r+1 \times r+3}{r \times r+4} b^2$, existente r
 æquali numero terminorum præcedentium quadrupli-
 cato: sic terminus sextus, quia tunc $r = 5 \times 4 = 20$,
 est

est æqualis termino quinto $\frac{2}{2} \dots \frac{17 \times 19}{16 \times 20} b^2$ ducto in $\frac{21 \times 23}{20 \times 24} b^2$. Quamobrem termino quovis hujus seriei

dicto B, terminus subsequens erit $Bb^2 \times \frac{r+1 \times r+3}{r \times r+4}$,

sive $Bb^2 \times 1 + \frac{3}{r \times r+4}$, et manente jam eodem valore

numeri r , termini reliqui erunt, $Bb^4 \times 1 + \frac{3}{r \times r+4}$

$\times 1 + \frac{3}{r+4 \times r+8}$, $Bb^6 \times 1 + \frac{3}{r \times r+4} \times 1 + \frac{3}{r+4 \times r+8}$

$\times 1 + \frac{3}{r+8 \times r+12}$, &c. Sed si fuerit r numerus ali-

quantum magnus, erit $1 + \frac{3}{r \times r+4} \times 1 + \frac{3}{r+4 \times r+8}$

$= 1 + \frac{3}{r \times r+4} + \frac{3}{r+4 \times r+8}$ quamproximè, et

$1 + \frac{3}{r \times r+4} \times 1 + \frac{3}{r+4 \times r+8} \times 1 + \frac{3}{r+8 \times r+12} =$

$1 + \frac{3}{r \times r+4} + \frac{3}{r+4 \times r+8} + \frac{3}{r+8 \times r+12}$, &c. Unde

termini omnes prædicti incipientes à termino B erunt

$$\begin{aligned}
 B + Bb^2 + Bb^4 + Bb^6 + Bb^8 + \dots &= \frac{B}{1-b^2} \\
 + \frac{3Bb^2}{r \times r+4} + \frac{3Bb^4}{r \times r+4} + \frac{3Bb^6}{r \times r+4} + \frac{3Bb^8}{r \times r+4} + \dots &= \frac{3B}{r \times r+4} \times \frac{b^2}{1-b^2} \\
 + \frac{3Bb^4}{r+4 \times r+8} + \frac{3Bb^6}{r+4 \times r+8} + \frac{3Bb^8}{r+4 \times r+8} + \dots &= \frac{3B}{r+4 \times r+8} \times \frac{b^4}{1-b^2} \\
 + \frac{3Bb^6}{r+8 \times r+12} + \frac{3Bb^8}{r+8 \times r+12} + \dots &= \frac{3B}{r+8 \times r+12} \times \frac{b^6}{1-b^2} \\
 + \frac{3Bb^8}{r+12 \times r+16} + \dots &= \frac{3B}{r+12 \times r+16} \times \frac{b^8}{1-b^2} \\
 \dots & \dots
 \end{aligned}$$

Ac

Ac proinde erit

$$S = \frac{1}{2}b + \frac{1}{2} \times \frac{5 \times 7}{4 \times 8} b^3 + \frac{1}{2} \times \frac{5 \times 7}{4 \times 8} \times \frac{9 \times 11}{8 \times 12} b^5 + \frac{1}{2} \times \frac{5 \times 7}{4 \times 8} \times \frac{9 \times 11}{8 \times 12} \times \frac{13 \times 15}{12 \times 16} b^7 +, \&c. + \frac{B}{1-b^2}$$

$$\times 1 + \frac{3b^2}{r \times r + 4} + \frac{2b^4}{r + 4 \times r + 8} + \frac{2b^6}{r + 4 \times r + 12} + \frac{3b}{r + 12 \times r + 16} +, \&c.$$

Itaque si, computatis, exempli gratiâ, quindecim terminis, decimus sextus designetur per B, erit $r = 15 \times 4 = 60$, et summa terminorum quindecim illorum addita

$$\text{summæ seriei } \frac{B}{1-b^2} \times 1 + \frac{3b^2}{r \times r + 4} + \frac{3b^4}{r + 4 \times r + 8} +, \&c.$$

dabit valorem coefficientis S.

Determinatis hoc pacto quantitibus assumptis R, S, T, &c. jam ut ad expressiones virium revertamur, vis

TR ad radium QS perpendicularis erat $\frac{\phi k}{z^3} - \frac{\phi}{k^2}$

$\times \sin. QSP$; sed posuimus angulum $QSP = \frac{1}{n}s$,

estque $\frac{1}{z^3} = \frac{1}{r^3} \sin R + S \cos. \frac{1}{n}s + T \cos. \frac{2}{n}s + V$

$\cos. \frac{3}{n}s + W \cos. \frac{4}{n}s +, \&c.$

$$\text{Unde vis } TR = \frac{\phi k}{r^3} \sin R - \frac{r^3}{k^3} - \frac{T}{2} \times \sin. \frac{1}{n}s$$

$$+ \frac{S-V}{2} \sin. \frac{2}{n}s + \frac{T-W}{2} \sin. \frac{3}{n}s + \frac{V-X}{2} \sin. \frac{4}{n}s$$

$$+, \&c.$$

Et vis quæ planetam Q distrahît à Sole in directione radii QS erat $\frac{\phi}{z^3} \times k \cos. QSP - 1 - \frac{\phi}{r^3} \cos. QSP$,

$$\text{hoc est, } \frac{\phi}{r^3} \sin \frac{kS}{2} - R + kR + \frac{kT}{2} - \frac{r^3}{k^2} - S$$

$$\times \cos.$$

$$\times \cos. \frac{1}{n} s + \frac{kS+kV-2T}{2} \times \cos. \frac{2}{n} s + \frac{kT+kW-2V}{2} \\ \cos. \frac{3}{n} s + \frac{kV+kX-2W}{3} \cos. \frac{4}{n} s +, \&c. \quad Q. E. I.$$

PROPOSITIO II. PROBLEMA.

Inæqualitates motûs planetæ interioris ex viribus prædictis ortas investigare.

Exeant simul planetæ P, Q (Fig. 2.) de locis D, C, ubi jacebant in eâdem rectâ cum Sole posito in S, et post aliquod temporis spatium reperiuntur in P et Q, et jungantur SP, SQ, PQ. Efto CS = 1, et arcus circularis CQ five angulus CSQ = s; denotent præterea P et Q respectivè tempora periodica planetarum P et Q, eritque ang. QSC : ang. PSD :: P : Q, adeoque angulus QSP : ang. QSC :: P - Q

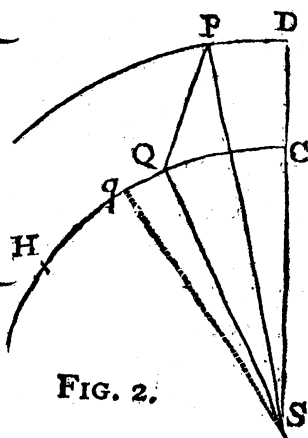


FIG. 2.

: P, unde ang. QSP = $\frac{1}{n} s$, posito $n = \frac{P}{P-Q}$.

Vis attractionis Solis ad distantiam QS, et tempus quo corpus, eâdem vi uniformiter agente, impulsûm acquirere posset eam velocitatem, qua planeta Q in circulo CQ revolvitur, tûm illa ipsa velocitas, exponantur figillatim per unitatem; et si, sumpto arcu CH = CS = 1, CH exprimat tempus illud unitati æquale, arcus quilibet quàm minimus Qq exprimet tempus quo uniformi illâ velocitate describitur.

Vol. LII.

Qq

Unde,

Unde, cum velocitates viribus quibufvis constantibus genitæ fint ut ipfæ vires et tempora, quibus hæ velocitates generantur, conjunctim ; erit velocitas 1 planetæ Q in circulo CQ revolvantis ad incrementum vel decrementum velocitatis vi Z genitum (fcripto nempe Z pro vi planetæ P normaliter ad radium QS agente, prout est in propofitione præcedente definita) quo tempore planeta Q describit arcum quàm minimum Qq, ut vis attractionis Solis 1 ducta in tempus CH five 1, ad vim Z ductam in tempus descriptionis arcûs Qq five in ipsum in arcum Qq: adeoque incrementum vel decrementum velocitatis vi Z genitum, quo tempore describitur arcus Qq, exprimeretur per $Z \times Qq$ five $Z \times s$.

Est autem $Z = \frac{\phi k}{t^3}$ in $R - \frac{t^3}{k^3} - \frac{T}{2} \times \sin. \frac{1}{n} s$
 $+ \frac{S-V}{2} \sin. \frac{2}{n} s + \frac{T-W}{2} \sin. \frac{3}{n} s +, \&c.$ et hac
 quantitate ductâ in s , tùm sumptâ fluente, prodit ve-
 locitatis acccleratio five retardatio, quam voco U,
 genita quo tempore describitur à planeta Q arcus CQ,
 æqualis $\frac{\phi k n}{t^3}$ in $R - \frac{t^3}{k^3} - \frac{T}{2} \times \sin. \text{vers.} \frac{1}{n} s + \frac{S-V}{4}$
 $\sin. \text{vers.} \frac{2}{n} s + \frac{T-W}{6} \sin. \text{vers.} \frac{3}{n} s + \frac{V-X}{8}$
 $\sin. \text{vers.} \frac{4}{n} s +, \&c.$ five pofito $b = R - \frac{t^3}{k^3} - \frac{T}{2}$
 $+ \frac{S-V}{4} + \frac{T-W}{6} + \frac{V-X}{8} +, \&c.$ $U = \frac{\phi k n}{t^3}$
 in $b - R - \frac{t^3}{k^3} - \frac{T}{2} \times \cos. \frac{1}{n} s - \frac{S-V}{4} \cos. \frac{2}{n} s$
 $- \frac{T-W}{6} \cos. \frac{3}{n} s - \frac{V-X}{8} \cos. \frac{4}{n} s -, \&c.$

Hoc pacto obtinetur variatio velocitatis in hypothesi quòd revolvatur planeta Q semper ad eandem distantiam à Sole, quod in præcedenti calculo supponi potest, cùm tantillùm varietur distantia SQ actione planetæ P .

Hoc facto, ut investigetur variatio distantiae planetæ Q à Sole, fingamus planetam descripsisse, non arcum circulem CQ , sed arcum curvæ Cr (Fig. 3.) et reperiri in puncto r ubi radius SQ productus fecat curvam.

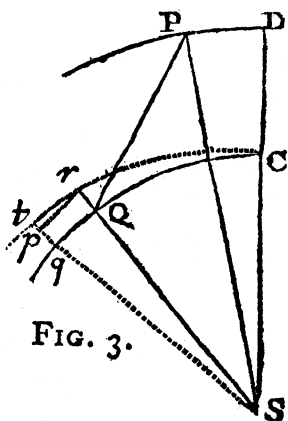


FIG. 3.

Ducatur recta St vicinissima ipsi SQ occurrens circulo et curvæ q et t ; tùm centro S et radio Sr describatur arcus rp , sitque $Sr = x$. Si planeta Q urgeretur solâ vi tendente ad centrum S , describeret areas temporibus proportionales, atque aded, cùm ipsius velocitas angularis in loco C supponatur esse 1 , in loco r foret

æqualis $\frac{1}{x}$; sed in illo quem exhibet schema situ planetarum minuitur hæc velocitas quantitate U suprâ

definitâ, unde velocitas angularis in loco r erit $\frac{1}{x} - U$;

et tempus, quo describeretur arcus Qq velocitate 1 , est ad tempus quo describitur arcus rp velocitate

$\frac{1}{x} - U$, ut Qq ad $\frac{rp}{\frac{1}{x} - U}$, hoc est, ut s ad $\frac{x s}{\frac{1}{x} - U}$;

unde, cùm s exprimat ex jam dictis tempus descriptionis arcûs Qq velocitate 1 , exprimet quantitas

titas $\frac{x^2}{\frac{1}{x} - U}$ tempus quo describitur arcus $r p$ velocitate

$\frac{1}{x} - U$. His positis, quoniam planetæ Q recessus à centro vel ad idem accessus pendet ex differentiâ virium, centrifugæ scilicet et centripetæ, quibus urge-
tur in Q ; si hæc differentia virium dicatur P , et v denotet velocitatem ascensûs vel descensûs planetæ Q secundum radium SQ , per idem planè ratiocinium, quod mox usurpavimus in investigatione velocitatis U , habebitur $v = P \times \frac{x^2}{\frac{1}{x} - U}$.

Quoniam ex hypothesi planeta Q , sepositâ actione planetæ P , describeret circulum, vires (centripeta et centrifuga) sibi invicem et unitati forent æquales: existente autem planetâ Q in r , ipsius attractio in Solem est $\frac{1}{k^2}$, ex qua auferenda est vis ea qua juxta propositionem præcedentem distrahitur à Sole, nimirum $\frac{0}{r^3}$ in $A + B \cos. \frac{1}{n} s + C \cos. \frac{2}{n} s + D \cos. \frac{3}{n} s + E \cos. \frac{4}{n} s +$, &c. positis $A = \frac{kS}{2} - R$, $B = kR + \frac{kT}{2} - \frac{r^3}{k^2} - S$, $C = \frac{kS + kV - 2T}{2}$, $D = \frac{kT + kW - 2V}{2}$, $E = \frac{kV + kX - 2W}{2}$, &c. atque harum virium differentia componit vim centripetam.

Vis autem centrifuga est semper in ratione duplicatâ areæ temporis momento descriptæ directè et triplicatâ distantie inversè; unde si hæc vis fuerit æqualis

1, ubi incepit planeta movere in C, erit æqualis
 $x^2 \times \frac{1}{x} - U^2 \times \frac{1}{x^3} = \frac{1}{x} \times \frac{1}{x} - U^2$ ubi movetur in r.

Differentia igitur inter vim centrifugam et centri-
 petam, qua urgetur planeta in r supra designata

per P, est $\frac{1}{x} \times \frac{1}{x} - U^2 - \frac{1}{x^2} + \frac{\phi}{x^3}$
 $\times A + B \cos. \frac{1}{n}s + C \cos. \frac{2}{n}s + D \cos. \frac{3}{n}s +, \&c.$

hincque habetur $\dot{v} = s \times \frac{1}{x} - U - \frac{s}{x \times \frac{1}{x} - U} + \frac{\phi}{x^3}$

$\times \frac{x s}{\frac{1}{x} - U} \times A + B \cos. \frac{1}{n}s + C \cos. \frac{2}{n}s +, \&c.$

Vires, quibus perturbatur motus planetæ Q, cùm
 exprimantur seriebus quorum termini ducuntur in
 finum vel cofinum anguli $\frac{1}{n}s$, vel anguli hujus mul-
 tiplicis, fingemus differentiam inter distantias S Q et
 S r exprimi serie fimili, ac propterea ponemus $x =$

$1 - Q + K \cos. \frac{1}{n}s + L \cos. \frac{2}{n}s + M \cos. \frac{3}{n}s$
 $+ N \cos. \frac{4}{n}s, \&c.$ existente $Q = K + L + M$
 $+ N +, \&c.$ ut fit S r, five $x = 1$, ubi planetæ
 Q et P incipiunt movere à lineâ conjunctionis S C D.
 Quantitates autem assumptæ K, L, M, &c. sunt exi-
 guæ, ideoque erit $\frac{1}{x} = 1 + Q - K \cos. \frac{1}{n}s - L$
 $\cos. \frac{2}{n}s - M \cos. \frac{3}{n}s - N \cos. \frac{4}{n}s -, \&c.$ quam-
 proximè.

proximè. Substituantur ergò in æquatione suprà traditâ valores quantitatum x , $\frac{1}{x}$, et U ; et sumptâ fluente, rejeclis iis terminis qui ducuntur in altiore quàm unam dimensionem quantitatum ϕ , Q , K ,

$$L, \text{ \&c. } \text{prodit } v = - \frac{2\phi k b n}{t^3} - \frac{\phi}{t^3} A - Q \times s$$

$$+ \frac{2\phi k n}{t^3} \times R - \frac{t^3}{k^3} - \frac{T}{2} + \frac{\phi}{t^3} B - K \times n \times \sin. \frac{1}{n} s$$

$$+ \frac{\phi k n}{t^3} \times \frac{S - V}{4} + \frac{\phi}{t^3} \times \frac{C}{2} - \frac{L}{2} \times n \times \sin. \frac{2}{n} s$$

$$+ \frac{\phi k n}{t^3} \times \frac{T - W}{9} + \frac{\phi}{t^3} \times \frac{D}{3} - \frac{M}{3} \times n \times \sin. \frac{3}{n} s$$

$$+ \frac{\phi k n}{t^3} \times \frac{V - X}{16} + \frac{\phi}{t^3} \times \frac{E}{4} - \frac{N}{4} \times n \times \sin. \frac{4}{n} s +, \text{ \&c.}$$

+ Z , designante Z quantitatem idoneam qua compleatur fluens. At, quoniam velocitas v supponitur nulla evadere, non solum ubi s , five arcus $CQ = 0$, id est, ubi planetæ versantur in primâ illâ conjunctione, sed etiam in omnibus aliis conjunctionibus subsequen-

tibus, hoc est, ubi est angulus $\frac{1}{n} s$, seu $PSQ = 0$,

vel $= r \times 180^\circ$, scripto scilicet r pro quovis ex numeris naturalibus 1, 2, 3, 4, &c. fiet $Z =$

$$\frac{2\phi k b n}{t^3} - \frac{\phi}{t^3} A - Q \times s \text{ adeoque}$$

$$v = \frac{2\phi k n}{t^3} \times R - \frac{t^3}{k^3} - \frac{T}{2} + \frac{\phi}{t^3} B - K \times n \times \sin. \frac{1}{n} s$$

$$+ \frac{\phi k n}{t^3} \times \frac{S - V}{4} + \frac{\phi}{t^3} \times \frac{C}{2} - \frac{L}{2} \times n \times \sin. \frac{2}{n} s$$

$$+ \frac{\phi k n}{t^3} \times \frac{T - W}{9} + \frac{\phi}{t^3} \times \frac{D}{3} - \frac{M}{3} \times n \times \sin. \frac{3}{n} s$$

+

$$+ \frac{\phi k n}{t^3} \times \frac{V - X}{16} + \frac{\phi}{t^3} \times \frac{E}{4} - \frac{N}{4} \times n \times \sin. \frac{4}{n} s \\ +, \&c.$$

Deinde, cùm fit tp , five \dot{x} ad rp , five xs , ut velocitas v qua describitur tp ad velocitatem $\frac{1}{x} - U$ qua describitur rp , erit $\dot{x} = v \times \frac{xs}{\frac{1}{x} - U}$, five, quia va-

lor velocitatis v componitur ex quantitativis exiguis, $\dot{x} = vs$ quamproximè, et $\frac{\dot{x}}{s} = v$. Verùm etiam æquatio assumpta $x = 1 - Q + K \cos. \frac{1}{n} s + L \cos. \frac{2}{n} s + M \cos. \frac{3}{n} s +, \&c.$ dat $\frac{\dot{x}}{s} = -K \times \frac{1}{n} \sin. \frac{1}{n} s - L \times \frac{2}{n} \sin. \frac{2}{n} s - M \times \frac{3}{n} \sin. \frac{3}{n} s - N \times \frac{4}{n} \sin. \frac{4}{n} s, \&c.$

Habitis igitur duobus velocitatis v valoribus, eorum termini homologi statuantur æquales, atque inde obtinebuntur quantitates assumptæ, nempe

$$K = \frac{\phi}{t^3} \times \frac{n^2}{n^2 - 1} \times \frac{2t^3}{k^2} \times \overline{n + \frac{1}{2}} - kT \times \overline{n - \frac{1}{2}} - S \\ L = \frac{\phi}{2t^3} \times \frac{n^2}{n^2 - 4} \times \overline{kS \times n + 1 - kV \times n - 1 - 2T} \\ M = \frac{\phi}{3t^3} \times \frac{n^2}{n^2 - 9} \times \overline{kT \times n + \frac{1}{2} - kW \times n - \frac{1}{2} - 3V} \\ N = \frac{\phi}{4t^3} \times \frac{n^2}{n^2 - 16} \times \overline{kV \times n + 2 - kX \times n - 2 - 4W} \\ \&c.$$

indeque manifesta fit harum quantitatum progressio :
atque

atque hoc pacto habetur semper distantia x planetæ Q à Sole.

Jam ut definiatur planetæ Q motus verus qui designatur per s , dicatur w motus medius, five, quod perinde est, tempus quo planeta descriperit arcum quemlibet Cr; atque ex demonstratis est $\dot{w} =$

$$\frac{x \dot{s}}{\frac{1}{x} - U}; \text{ unde, substitutis valoribus quantitatum, } x,$$

$\frac{1}{x}$, et U, et sumptâ fluente, emergit

$$w = 1 - 2Q + \frac{\phi k b n}{3} \times s + 2nK - \frac{\phi k n^2}{t^3} \times R - \frac{t^3}{k^3} - \frac{T}{2}$$

$$\times \text{fin. } \frac{1}{n}s + nL - \frac{\phi k n^2}{8t^3} \times S - \bar{V} \times \text{fin. } \frac{2}{n}s$$

$$+ \frac{2nM}{3} - \frac{\phi k n^2}{18t^3} \times T - \bar{W} \times \text{fin. } \frac{3}{n}s$$

$$+ \frac{nN}{2} - \frac{\phi k n^2}{32t^3} \times \bar{V} - \bar{X} \times \text{fin. } \frac{4}{n}s +, \&c. + Z$$

denotante Z quantitatem idoneam ut compleatur fluens. Sed, quia motus verus medio æqualis evadere supponitur in qualibet planetarum P et Q conjunctione

cum Sole, id est, ubi angulus PSQ five $\frac{1}{n}s$ æquatur,

vel nihilo, vel angulo $r \times 180^\circ$, exhibente r quemvis ex numeris naturalibus 1, 2, 3, 4, &c. erit Z =

$$2Q - \frac{\phi k b n}{t^3} \times s. \text{ Ponantur igitur } F = -2nK + \frac{\phi k n^2}{t^3}$$

$$\times R - \frac{t^3}{k^3} - \frac{T}{2}, G = -nL + \frac{\phi k n^2}{8t^3} \times S - \bar{V},$$

$$H = -\frac{2nM}{3} + \frac{\phi k n^2}{18t^3} \times T - \bar{W}, I = -\frac{nN}{2} + \frac{\phi k n^2}{32t^3}$$

X

+ V — X, &c. eritque motus verus, five $s = w$
 + F fin. $\frac{1}{n}s$ + G x fin. $\frac{2}{n}s$ + H x fin. $\frac{3}{n}s$ + I
 x fin. $\frac{1}{n}s$ +, &c. vel, quia parum admodum differt
 motus verus à motu medio $s = w$ + F x fin. $\frac{1}{n}w$
 + G x fin. $\frac{2}{n}w$ + H x fin. $\frac{3}{n}w$ + I x fin. $\frac{4}{n}w$ +,
 &c. Q. E. I.

COROLL. I.

His ita generatim definitis, ut specialis eliciatur in
 motu cujuspian planetæ inæqualitatum mensura, de-
 terminandæ sunt quantitates assumptæ.

Itaque planeta P designet Terram, planeta Q Ve-
 nerem, et quoniam est distantia Terræ ad distantiam
 Veneris à Sole ut 100000 ad 72333, hæc erit ratio
 k ad 1, adeoque $k = \frac{100000}{72333}$, $kk + 1 = tt =$
 2.91129 , $b = \frac{2k}{t^3} = 0.94975$; atque inde per me-
 thodum in Prop. I^a. expositam prodibunt

R = 9.3925	V = 11.1964	Y = 5.3380
S = 16.6782	W = 8.8504	Z = 4.1029
T = 13.8877	X = 6.9045	&c.

Tum, existente periodo Terræ annuâ dierum
 365.2565, et periodo Veneris dierum 224.701, est
 ex jam dictis $n = \frac{365.2565}{365.2565 - 224.701} = 2.59866$;
 et cùm gravitas in Solem sit juxta Newtonum ad gra-
 vitatem in Terram, paribus distantis, ut 1 ad $\frac{1}{169444}$,
 erit $\phi = \frac{1}{169444}$.

Unde, redactis in numeros formulis in hac propositione datis, emergunt

$$\begin{array}{ll} K = 0.0000103 & N = - 0.0000065 \\ L = 0.0000444 & O = - 0.0000024 \\ M = 0.0000377 & O' = - 0.0000011, \text{ \&c.} \end{array}$$

Atque ex his tandem deducuntur

$$\begin{array}{ll} F = - 0.0000473 & I = 0.0000100 \\ G = - 0.0001078 & I' = 0.0000033 \\ H = - 0.0000684 & \text{ \&c.} \end{array}$$

Hinc ergo habentur valores coefficientium æquationis $s = w + F \times \sin. \frac{1}{n} w + G \times \sin. \frac{2}{n} w + H \times \sin. \frac{3}{n} w +$, &c. ubi s denotat motum Veneris verum, w motum medium, et $\frac{1}{n} w$ angulum PSQ five differentiam longitudinum heliocentricarum Terræ et Veneris; vel, reductis quantitibus F, G, H, &c. ad exprimendas more astronomico circuli partes, fit $s = w - 9''.76 \times \sin. \frac{1}{n} w - 22''.24 \times \sin. \frac{2}{n} w + 14''.11 \times \sin. \frac{3}{n} w + 2''.06 \times \sin. \frac{4}{n} w + 0''.68 \times \sin. \frac{5}{n} w +$, &c.

Ut exemplum apponam, esto angulus PSQ five $\frac{1}{n} w = 40^\circ$, ac prodibit $s = w - 15''.5$; motus igitur medius superat verum, eorumque differentia est $15''.5$.

Computatâ hoc pacto differentiâ inter motum Veneris verum et medium respectu Solis, sequenti modo innotescet quanta evadat cum e Terrâ spectatur. Esto

PSQ

Elongatio Ven. à Sole.	Correctio.	Elongatio Ven. à Sole.	Correctio.
° / "	"	° / "	"
0	0	46 19 50	0
5	0	46	+ 2.3
10	0	45	5.1
15	0	40	9.5
20	— 0.5	35	7.3
25	0.8	30	1.8
30	1.5	25	— 4.4
35	2.8	20	9.2
40	2.9	15	11.2
45	2.7	10	10.2
46	1.7	5	6.0
46 19 50	0	0	0

Exempli gratiâ, si Venus à conjunctione inferiore digressa motu suo medio discefferit à Sole angulo elongationis 40° , erit vera Veneris elongatio $40^{\circ} - 2''.9 = 39^{\circ} 59' 57''.1$: pariter, si ulteriùs delata Venus pervenerit ad eandem elongationem 40° , erit tunc vera Veneris elongatio $40^{\circ} 0' 9''.5$. Eadem omninò sunt correctiones et cum iisdem signis adhibendæ, ubi post conjunctionem superiorem eadem eveniunt elongationes.

COROLL. II.

Ex præcedentibus etiâ deducitur distantia Veneris à Sole pro quolibet ejus cum Terrâ et Sole aspectu, in

in hypothesi quod, seclusâ Terræ attractione, in orbitâ circulari revolveret. Sic, si angulus $\frac{1}{n}s$, seu PSQ fit 90° , vel 270° , æquatio $x = 1 - Q + K \cos. \frac{1}{n}s + L \cos. \frac{2}{n}s + M \cos. \frac{3}{n}s + N \cos. \frac{4}{n}s +$, &c. fit $x = 0.9999437$ circiter; et si fit PSQ = 180° , fit $x = 1.0000607$.

Unde, si distantia Veneris à Sole in conjunctione inferiore ponatur	- - -	} 10000000
In quadraturis cum Terrâ erit ipsius distantia	- - - - -	
	- - - - -	} 9999437
In conjunctione superiore erit	- - -	
	- - -	10000607

Item innotescit differentia inter tempus periodicum Veneris, quale nunc est, et tempus illud periodicum, quale foret, si unicâ Solis attractione in orbe circulari moveretur. Siquidem, cùm Venus post discessum suum à conjunctione ad eamdem redierit, æquatio generalis in propositione tradita, quæ exprimit relationem inter motum Veneris verum et medium, evadit

$$w = 1 - 2Q + \frac{\phi k h n}{i^3} \times s, \text{ five } w = 1.0000066 \times s$$

circiter: unde tempus periodicum Veneris est ad tempus illud alterum periodicum, ut 1.0000066 ad 1; adeoque, si nulla foret gravitatio Veneris in Terram, revolutionem suam circa Solem minutis duobus horæ primis citiùs perageret.

PROPOSITIO III. PROBLEMA.

In systemate duorum planetarum in orbitis circularibus circa Solem revolventium, motum nodorum orbitæ planetæ interioris, quatenus ex vi planetæ exterioris oritur, investigare.

Per motum nodorum hîc intelligendus est motus interfectionis plani orbis planetæ interioris cum plano orbis planetæ exterioris spectato ut immoto. Itaque esto Sol in S (Fig. 5.) et centro S atque radio SQ de-

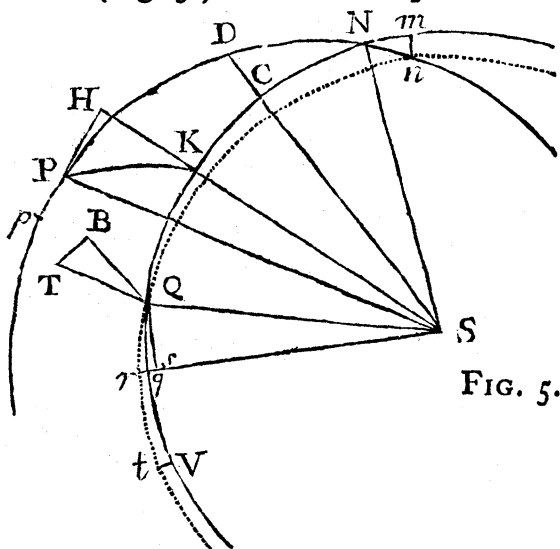


FIG. 5.

scribantur in superficie sphæræ duo circuli QN, PN, sese interfecantes in N, quorum prior QN designet situm plani orbis planetæ interioris Q, et posterior PN situm plani orbis planetæ exterioris, cujus locus sit in rectâ SP productâ. Eodem centro S et radio SP describatur circulus PK, cujus planum sit plano
SQN

SQN perpendicularare, secetque circulum QN in K ,
 et in SK demittatur perpendiculum PH : tum ductâ
 QT parallelâ rectæ SP et TB in planum SQN nor-
 mali, si linea QT exhibeat vim qua trahitur planeta
 Q in directione QT , seu SP , TB exhibebit vim qua
 distrahitur perpendiculariter à plano suæ orbitæ; erit-
 que triangulum QTB simile triangulo SPH , atque
 adeò, $TB : QT :: PH : SP :: \text{fin. } PK : 1$; deinde
 in triangulo sphærico rectangulo PKN habetur,
 $1 : \text{fin. } PN :: \text{fin. } PNK : \text{fin. } PK$; unde, conjunctis
 rationibus, et scripto c pro sinu anguli PNK ad ra-
 dium 1 , hoc est, pro sinu inclinationis orbis QN ad
 orbem PN , provenit $TB = QT \times c \times \text{fin. } PN$.
 Sumatur jam arcus quàm minimus Qq , ad quem
 erigitur lineola perpendicularis qr , æqualis duplo spatio
 quod planeta Q percurrere posset impellente vi TB
 quo tempore in orbe suo describeret arcum illum Qq ,
 et centro S descriptus circulus rQn secans circulum
 PN in n exhibebit situm orbis planetæ Q post tem-
 pus illud, nodo N translato in n ; atque in QN de-
 missò perpendiculo nm , et in Sq perpendiculo Qs ,
 erit angulus qQr , five NQn ad duplum angulum
 qQs , id est, ad angulum QSq , ut vis TB ad gra-
 vitatem (nempe 1) planetæ Q in Solem; hoc est,
 $\frac{nm}{\text{fin. } QN} : Qq :: TB : 1$; in triangulo autem rectan-
 gulo Nmn , est $Nn : nm :: 1 : c$; quare conjunctis his
 rationibus, prodit $Nn = \frac{TB \times \text{fin. } QN \times Qq}{c}$, sed
 suprâ invenimus $TB = QT \times c \times \text{fin. } PN$, unde fit
 $Nn = QT \times \text{fin. } PN \times \text{fin. } QN \times Qq$.
 Esto SC linea conjunctionis planetarum, fiatque, ut
 in propositione præcedente, arcus $CQ = s$, $Qq = s$,
 SQ

$SQ = 1$; et, quia inclinatio orbis QN ad orbem PN exigua supponitur, erit etiam hîc ang. $PSQ = \frac{1}{n}s$ quamproximè ; proindeque, posito arcu $CN = a$, erit $QN = s + a$ et $PN = s - \frac{1}{n}s + a$ quamproximè.

Porro, cùm lentissimè moveantur nodi, arcus CN spectari potest quasi invariabilis per multarum planetæ Q revolutionum seriem, atque adedò fluxio arcus QN eadem erit cum fluxione arcûs QC . His positis, ha-

bebitur $\sin. PN \times \sin. QN = \frac{1}{2} \cos. \frac{1}{n}s - \frac{1}{2} \cos. 2s - \frac{1}{n}s + 2a$, estque per propositionem primam

$QT = \frac{\phi^k}{z^3} - \frac{\phi}{k^2} = \frac{\phi^k}{t^3} \text{ in } R - \frac{t^3}{k^3} + S \cos. \frac{1}{n}s + T \cos. \frac{2}{n}s + V \cos. \frac{3}{n}s + W \cos. \frac{4}{n}s +, \&c.$ unde

substitutis his valoribus in æquatione $Nn = QT \times \sin. PN \times \sin. QN \times Qq$, et sumptâ fluente per methodum in Coroll. IV. lemmatis edoctam, prodibit summa omnium Nn , sive motus nodi, quo tempore planeta Q à loco conjunctionis C procedens

in orbe suo descriperit arcum CQ , æqualis $\frac{\phi^k n}{2 t^3}$ in

$\frac{S}{2n}s + R - \frac{t^3}{k^3} + \frac{T}{2} \times \sin. \frac{1}{n}s + \frac{S+V}{4} \sin. \frac{2}{n}s + \frac{T+W}{6} \sin. \frac{3}{n}s +, \&c. + \frac{\phi^k n}{2 t^3} \text{ in } Z \times \sin. 2a$

$- R - \frac{t^3}{k^3} \times \frac{1}{2n-1} \sin. 2s - \frac{1}{n}s + 2a - \frac{S}{2} \times \frac{1}{2n} \sin. 2s + 2a - \frac{S}{2} \times \frac{1}{2n-2} \sin. 2s - \frac{2}{n}s + 2a - \frac{T}{2}$

\times

$$\begin{aligned} & \times \frac{1}{2n-3} \sin. 2s - \frac{3}{n} s + 2a - \frac{T}{2} \times \frac{1}{2n+1} \sin. \\ & 2s + \frac{1}{n} s + 2a - \frac{V}{2} \times \frac{1}{2n-4} \sin. 2s - \frac{4}{n} s + 2a \\ & - \frac{V}{2} \times \frac{1}{2n+2} \sin. 2s + \frac{2}{n} s + 2a, \&c. \text{ existente} \\ Z = \frac{1}{2n-1} \sin R - \frac{t^3}{k^3} \times \frac{1}{(2n-1)^2} + \frac{S}{2n-2+2n} \\ & + \frac{T}{2n-3 \times 2n+1} + \frac{V}{2n-4 \times 2n+2} + \frac{W}{2n-5 \times 2n+3} \\ & +, \&c. \text{ atque in his seriebus patet terminorum pro-} \\ & \text{gressio. Q. E. I.} \end{aligned}$$

COROLL. I.

Hic liquet multas oriri in motu nodorum æquationes; sed quia minutæ sunt, et locum planetæ Q ferè nihil mutant, idèò satis erit rationem habere motûs nodorum medii et æquationis solûs *periodicæ*, qui sic ex præcedentibus deducuntur. Cum in planis parùm ad se inclinatis moveri supponantur planetæ P et Q, quoties revertentur ad conjunctionem, angulus PSQ,

sive $\frac{1}{n} s$, qui metitur eorum distantiam à se invicem, evadet $= 360^\circ$ vel $= r \times 360^\circ$, existente r numero integro: et quia, sumpto arcu quolibet A, est semper $\sin. r \times 360^\circ + A = \sin. A$; hinc, si computatur motus nodi pro tempore conjunctionum, expressio illa generalis et prolixa in propositione tradita in hanc

simplicem abit $\frac{\phi k}{2t^3} \times \frac{S}{2} s - nZ \times \sin. 2s + 2a - \sin. 2a$, sive per Coroll. I. lemmatis

$$\frac{\phi k}{2t^3} \times \frac{S}{2} s - 2nZ \times \sin. s \times \cos. s + 2a.$$

Hic est igitur motus nodorum factus, quo tempore planetæ P et Q à conjunctione provecti post quotlibet-
cunque revolutiones ad conjunctionem quamvis aliam
pervenerint, exhibente s arcum à planetâ Q in suâ or-
bitâ intereà descriptum. Terminus $\frac{\phi k}{2t^3} \times 2nZ \times \sin. s$
 $\times \cos. \overline{s + 2a}$ indicat æquationem *periodicam* et fa-
cillimè computatur: cùmque hæc æquatio modò fit
additiva, modò subtractiva, patet termino altero $\frac{\phi k}{2t^3} \times \frac{S}{2} s$
exprimi generatim motum nodi medium.

COROLL. II.

Esto planeta P Terra, Q Venus, et revolutionem
Veneris ab unâ conjunctione inferiore cum Terrâ ad
alteram vocemus, brevitatis gratiâ, revolutionem *syn-*
odicam; eritque post unam revolutionem synodicam
 $\frac{1}{n}s = 360^\circ$, proindeque $s = n \times 360^\circ = 935^\circ 31'$;
hic igitur est arcus descriptus à Venere inter duas
ejusdem generis conjunctiones. Hinc motus nodi
medius tempore revolutionis unius synodicæ, qui juxta
corollarium præcedens est $\frac{\phi k S}{4t^3} s$ fit $\frac{\phi k n S}{4t^3} = 360^\circ =$
 $23''.087$; atque hic motus imminutus in ratione tem-
poris periodici Terræ circa Solem ad revolutionem
Veneris synodicam, id est, in ratione 1 ad $n - 1$,
evadit $14''.44$, motus scilicet annuus nodorum Ve-
neris regressivus, qui spatio centum annorum fit
 $24' 4''$.

Æquatio periodica $\frac{\phi k n Z}{t^3} \times \sin. s \times \cos. \overline{s + 2a}$ ut
adhuc simplicior evadat, ponamus arcum a five CN
perexiguum

perexiguum esse vel nullum, id est, supponamus conjunctionem Terræ et Veneris fieri proximè in nodo, quemadmodum contingit hoc anno 1761, eritque

$$\text{æquatio periodica } \frac{\phi kn Z}{t^3} \times \text{fin. } s \times \text{cos. } s = \frac{\phi kn Z}{2 t^3}$$

$\times \text{fin. } 2s$. Cùm igitur fit $Z = 32.33$ circiter, for-

$$\text{mula } \frac{\phi ks}{4 t^3} s - \frac{\phi kn Z}{2 t^3} \text{fin. } 2s, \text{ quæ per corollarium præ-$$

cedens exprimit generatim motum nodi in quali-

bet serie revolutionum synodicarum confectum, fit

$$0.000006855 \times s - 14''.2 \times \text{fin. } 2s. \text{ Æquatio igitur}$$

periodica $14''.2 \times \text{fin. } 2s$, quam *generalem* voco, est

ut sinus dupli arcûs à Venere descripti in datâ serie re-

volutionum synodicarum, nec ultra $14''.2$ ascendit.

Jam, si pro s substituatur $935^\circ 31'$, erit $\text{fin. } 2s =$

$\text{fin. } 71^\circ 2'$, et regredientur nodi, in primâ revolutione

synodicâ post conjunctionem factam in nodo, per ar-

cum $23'' - 14''.2 \times \text{fin. } 71^\circ 2' = 10''$: et, si r de-

notet numerum quemcumque revolutionum synodi-

carum, motus nodi, peractis illis revolutionibus, erit

$r \times 23' - 14''.2 \times \text{fin. } r \times 71^\circ 2'$; pariterque, per-

actis revolutionibus quarum numerus est $r - 1$, idem

motus erit $r - 1 \times 23'' - 14''.2 \times \text{fin. } r - 1$

$\times 71^\circ 2'$; posterior motus ex priore auferatur, et re-

manebit $23'' - 14''.2 \times \text{fin. } r \times 71^\circ 2' - \text{fin. } r - 1 \times 71^\circ 2'$

$= 23'' - 14''.2 \times 2 \text{fin. } 35^\circ 31' \times \text{cos. } r \times 71^\circ 2' - 35^\circ 31'$

$= 23'' - 16''.5 \times \text{cos. } 2r - 1 \times 35^\circ 31'$ pro motu

nodi factò, tempore illius revolutionis synodicæ, cujus

locum in serie revolutionum indicat numerus r .

Exempli gratiâ, si desideretur motus nodi tempore

revolutionis quartæ synodicæ post conjunctionem

factam in nodo, erit $r = 4$, et regressus nodi erit

$23'' - 16''.5 \times \cos. 7 \times 35^\circ 31' = 29''$. Sic ope hujus formulæ $23'' - 16''.5 \times \cos. 2r - 1 \times 35^\circ 31'$ facilè computatur sequens tabula, quæ exhibet regressum nodi Veneris in plano eclipticæ, pro duodecim figillatim revolutionibus synodicis quæ proximè sequuntur conjunctionem Terræ et Veneris factam in nodo vel proximè ad nodum.

In revol. Ven. synod.	Regressus nodi Ven.	In revol. Ven. synod.	Regressus nodi Ven.
	//		//
1 ^a .	10	7 ^a .	26
2 ^a .	28	8 ^a .	39
3 ^a .	39	9 ^a .	30
4 ^a .	29	10 ^a .	11
5 ^a .	10	11 ^a .	8
6 ^a .	9	12 ^a .	25

Qui motus potest, cùm libuerit, ad annos communes reduci.

Denique patet æquationem periodicam, nempe $16''.5 \times \cos. 2r - 1 \times 35^\circ 31'$, quam *specialem* appello, ubi maxima est, evadere $16''\frac{1}{2}$; ac proinde regressum nodi in unâ revolutione synodicâ nusquam superare $39''\frac{1}{2}$, nec minorem esse $6''\frac{1}{2}$.

PROPOSITIO IV. PROBLEMA.

Iisdem positis, variationem inclinationis orbis planetæ interioris ad planum orbis planetæ exterioris determinare.

Esto NQV (Fig. 5.) quadrans circuli, cui erigatur perpendicularis Vt occurrens arcui nQr producto in t, eritque Vt mensura variationis inclinationis orbis NQV factæ quo tempore nodus N transfertur in n. Est autem Vt : nm :: fin. QV five cof. QN : fin. QN, atque nm : Nn :: c : 1, c denotante finum inclinationis orbis QN ad orbem PN, adeoque Vt : Nn :: c

× cof. QN : fin. QN; unde $Vt = Nn \times \frac{c \times \text{cof. QN}}{\text{fin. QN}}$,

five, quia per propositionem superiorem habetur Nn = QT × fin. PN × fin. QN × Qq, $Vt = c \times QT \times \text{fin. PN} \times \text{cof. QN} \times Qq$. Hinc, cum sit fin. PN

× cof. QN = $\frac{1}{2} \text{fin. } 2s - \frac{1}{n}s + 2a - \frac{1}{2} \text{fin. } \frac{1}{n}s$,

sumptâ fluente prodit variatio inclinationis, quo tempore planeta Q à loco conjunctionis C movetur per arcum CQ, æqualis

$-\frac{\phi c k n}{2t^3} \text{ in } R - \frac{t^3}{k^3} - \frac{T}{2} \times \text{fin. verf.}$

$\frac{1}{n}s + \frac{S-V}{4} \text{ fin. verf. } \frac{2}{n}s + \frac{T-W}{6} \text{ fin. verf. } \frac{3}{n}s$

$+ \frac{V-X}{8} \text{ fin. verf. } \frac{4}{n}s +, \&c. + \frac{\phi c k n}{2t^3} \text{ in } -Z$

$\times \text{fin. verf. } 2a + R - \frac{t^3}{k^3} \times \frac{1}{2n-1} \text{ fin. verf. } 2s - \frac{1}{n}s + 2a$

$+ \frac{S}{2} \times \frac{1}{2n} \text{ fin. verf. } 2s + 2a + \frac{S}{2} \times \frac{1}{2n-2} \text{ fin. verf.}$

$$\begin{aligned}
 & \overline{2s - \frac{2}{n}s + 2a + \frac{T}{2} \times \frac{1}{2n+1} \text{ fin. verf. } 2s + \frac{1}{n}s + 2a} \\
 & + \frac{T}{2} \times \frac{1}{2n-3} \text{ fin. verf. } 2s - \frac{3}{n}s + 2a + \frac{V}{2} \times \frac{1}{2n+2} \\
 & \text{fin. verf. } 2s + \frac{2}{n}s + 2a + \frac{V}{2} \times \frac{1}{2n-4} \text{ fin. verf.} \\
 & \overline{2s - \frac{4}{n}s + 2a + \frac{W}{2} \times \frac{1}{2n+3} \text{ fin. verf. } 2s + \frac{3}{n}s + 2a} \\
 & + \frac{W}{2} \times \frac{1}{2n-5} \text{ fin. verf. } 2s - \frac{5}{n}s + 2a, \text{ \&c.}
 \end{aligned}$$

Existente hîc eodem valore quantitatis Z ac in propositione præcedente. Q. E. I.

C O R O L L.

Si computetur variatio inclinationis pro tempore conjunctionum, facîle obtinebitur; hæc enim per formulam in propositione traditam evadit $\frac{\phi c k n}{2t^3} \times Z$
 $\times \text{fin. verf. } 2s + 2a - \text{fin. verf. } 2a$ quæ itèm, si
 prima conjunctionum, à qua sumitur motûs exordium,
 statuatur in nodo, fit $\frac{\phi c k n}{1t^3} \times Z \times \text{fin. verf. } 2s$.

Hoc est igitur decrementum inclinationis orbis planetæ Q factum in qualibet serie revolutionum ad conjunctionem, designante s arcum intereà à planetâ circa Solem descriptum. Conferatur hæc inclinationis variatio cum æquatione nodi periodicâ eodem tempore genitâ, prout in propositione superiore definitur, et patebit priorem esse ad posteriorem ut $c \times \text{fin. verf. } 2s$ ad $\text{fin. } 2s$.

Ut ad orbem Veneris hæc transferantur, quem si inclinari ad orbem Terræ supponatur angulo $3^\circ 23' 20''$,
 erit

erit $\frac{\phi c k n}{2 t^3} \times Z \times \sin. \text{verf. } 2s = 0''.84 \times \sin. \text{verf. } 2s.$

Unde palàm fit: 1°. in quacumque serie revolutionum synodicarum, post conjunctionem factam in nodo, decrementum inclinationis orbitæ Veneris ad eclipticam non superare $2 \times 0''.84 = 1''.68$, quod è Terrâ spectatum evadit $4''.4$: 2°. cùm, peractâ unâ revolutione synodicâ, fit $\sin. \text{verf. } 2s = \sin. \text{verf. } 71^\circ 2'$, inclinationis decrementum pro qualibet serie revolutionum synodicarum quarum numerus est r , effe $0''.84 \times \sin. \text{verf. } r \times 71^\circ 2'$, et pro serie revolutionum quarum numerus est $r - 1$, effe $0''.84 \times \sin. \text{verf. } r - 1 \times 71^\circ 2'$; unde horum decrementorum differentia

$0''.84 \times \sin. \text{verf. } r \times 71^\circ 2' - \sin. \text{verf. } r - 1 \times 71^\circ 2' = 0''.84 \times 2 \sin. 35^\circ 31' \times \sin. 2r - 1 \times 35^\circ 31' = 0''.98 \times \sin. 2r - 1 \times 35^\circ 31'$, exprimit variationem inclinationis genitam tempore revolutionis synodicæ illius, cujus locum in serie revolutionum denotat numerus r : atque hæc variatio, ut patet, nusquam excedit $0''.98$ è Sole conspecta, quæ spectatori in centro Terræ collocato sub angulo $2''\frac{1}{2}$ apparebit. Cum igitur tantilla sit orbitæ Veneris inclinationis variatio, non videtur operæ pretium de eâ ulterius exquirere.

Demonstratis, quæ ad perturbationem motûs planetæ interioris spectant, superest ut, quibus perturbationibus afficiatur motus planetæ exterioris, vicissim expendamus.

PROPOSITIO V. PROBLEMA.

In systemate duorum planetarum circa Solem in orbitibus penè circularibus revolventium, determinare vim planetæ interioris ad perturbandum motum exterioris.

Simili ratiocinio ei, quod in propositione primâ usurpavimus, etiam hoc problema solvitur. Itaque posita unitate pro distantia planetæ P à Sole, ubi ambo planetæ P et Q conjunguntur cum Sole, (Fig. 1.) fiat $SP = x$, $SQ = k$, $PQ = z$. Sit 1 ad ϕ ut gravitatio planetæ P in Solem in distantia 1 ad ejusdem planetæ P gravitationem in planetam Q in eadem distantia, eritque $\frac{\phi}{z^2}$ gravitas planetæ P in planetam Q in distantia PQ. Productâ, si opus est, PQ ad O ut fit $PO = \frac{\phi}{z^2}$, et ductâ OI parallelâ rectæ QS occurrente PS productæ in I, resolvatur vis PO in vires PI et OI, eritque propter similia triangula PQS, POI, vis OI $= \frac{PO \times QS}{PQ} = \frac{\phi k}{z^3}$, atque vis PI $= \frac{PO \times PS}{PQ} = \frac{\phi x}{z^3}$ five vis PI $= \frac{\phi}{z^3}$ quamproximè. Vis OI impellit planetam P in directione parallelâ rectæ SQ, et in eundem sensum urgetur Sol vi $\frac{\phi}{k^2}$ qua gravitat in planetam Q: excessu igitur solo vis prioris supra posteriorem, nempe $\frac{\phi k}{z^3} - \frac{\phi}{k^2}$, censendus est urgeri planeta P in directione parallelâ rectæ SQ.

Porro

Porro vis $\frac{\phi k}{z^3} - \frac{\phi}{k^2}$ ea pars, quæ agit perpendiculariter ad radium PS, est $\frac{\phi k}{z^3} - \frac{\phi}{k^2} \times \sin. PSQ$, atque altera pars, quæ amovet planetam P à Sole secundum PS, est $\frac{\phi k}{z^3} - \frac{\phi}{k^2} \times \cos. PSQ$. Auferatur hæc posterior vis ex vi PI, et manebit vis $\frac{\phi}{z^3} + \frac{\phi}{k^2} - \frac{\phi k}{z^3} \times \cos. PSQ$, qua planeta P urgetur versus Solem.

Est DCS (Fig. 2.) linea conjunctionis planetarum, et arcus DP, five angulus DSP vocetur s , denotentque P et Q respectivè tempora periodica planetarum P et Q, eritque, posito $n = \frac{Q}{P - Q}$, ang. PSQ = $\frac{1}{n}s$. Tum, si fiat $t^2 = 1 + kk$, et $b = \frac{2k}{t^2}$, erit uti in Prop. I. exposuimus, $z^2 = t^2 \times 1 - b \cos. \frac{1}{n}s$, atque $\frac{1}{z^3} = \frac{1}{t^3} \times R + S \cos. \frac{1}{n}s + T \cos. \frac{2}{n}s + V \times \cos. \frac{3}{n}s + \&c.$ et quemadmodum ibi erat $b = \frac{2PS \times SQ}{PS^2 + SQ^2}$, hîc item est $b = \frac{2PS \times SQ}{PS^2 + SQ^2}$, adeoque valores quantitatum assumptarum R, S, T, &c. iidem hîc sunt ac in propositione primâ.

Unde vis $\frac{\phi k}{z^3} - \frac{\phi}{k^2} \times \sin. PSQ$, qua sollicitatur planeta P in directione ad radium PS perpendiculari, sic exprimetur $\frac{\phi k}{t^3}$ in $R - \frac{t^3}{k^3} - \frac{T}{2} \times \sin. \frac{1}{n}s + \frac{S - V}{2} \sin.$

$$\text{fin. } \frac{2}{n} s + \frac{T - W}{2} \text{ fin. } \frac{3}{n} s + \frac{V - X}{2} \text{ fin. } \frac{4}{n} s +, \&c.$$

Et vis $\frac{\phi}{z^3} + \frac{\phi}{k^2} - \frac{\phi k}{z^3} \times \text{cof. PSQ}$, qua urgetur planeta P in Solem secundum radium PS, fiet

$$\frac{\phi}{t^3} \text{ in } R - \frac{kS}{2} - kR + \frac{kT}{2} - \frac{t^3}{k^2} - S \times \text{cof. } \frac{1}{n} s$$

$$- \frac{kS + kV - 2T}{2} \text{ cof. } \frac{2}{n} s - \frac{kT + kW - 2V}{2} \text{ cof. } \frac{3}{n} s$$

$$- \frac{kV + kX - 2W}{2} \text{ cof. } \frac{4}{n} s +, \&c. \quad \text{Q. E. I.}$$

PROPOSITIO VI. PROBLEMA.

Inæqualitates motûs planetæ exterioris ex viribus prædictis ortas investigare.

Per analysim in propositione secundâ institutam vis ad radium PS perpendicularis generabit accelerationem, vel retardationem velocitatis, dum arcus quilibet DP describitur à planeta P, æqualem $\frac{\phi k n}{t^3}$

$$\text{in } b = R - \frac{t^3}{k^3} - \frac{T}{2} \times \text{cof. } \frac{1}{n} s - \frac{S - V}{4} \text{ cof. } \frac{2}{n} s$$

$$- \frac{T - W}{6} \text{ cof. } \frac{3}{n} s - \frac{V - X}{8} \text{ cof. } \frac{4}{n} s -, \&c. = U$$

existente $b = R - \frac{t^3}{k^3} - \frac{T}{2} + \frac{S - V}{4} + \frac{T - W}{6}$

$$+ \frac{V - X}{8} +, \&c.$$

Deinde si scribatur p pro vi illâ planetæ Q qua urgetur planeta P in Solem, prout in propositione præcedente definita est, et v pro velocitate ascensûs vel descensûs planetæ P secundum radium PS, et jam

supponatur $SP = x = 1 - Q + K \cos. \frac{1}{n}s + L \cos. \frac{2}{n}s + M \cos. \frac{3}{n}s + N \cos. \frac{4}{n}s +, \&c.$ existente $Q = K + L + M + N +, \&c.$ erit $\frac{1}{x^2} + p$ vis centripeta planetæ P, et $\frac{1}{x} \times \sqrt{\frac{1}{x} + U}$ ejusdem vis centrifuga, atque inde habebitur $\dot{v} = \frac{\frac{1}{x^2} + p - \frac{1}{x} \times \sqrt{\frac{1}{x} + U}}{\frac{1}{x} + U} \times \frac{x \dot{s}}{\frac{1}{x} + U}.$

Tum restitutis valoribus quantitatum U, p, x, et prosequendo calculum prout in Prop. II. positis

$$A = Kn + \frac{2\phi kn^2}{t^3} \times R - \frac{t^3 - T}{k^3 - \frac{T}{2}} - \frac{\phi n}{t^3} \times kR - \frac{t^3 - S}{k^2 - \frac{T}{2}} + \frac{kT}{2}$$

$$B = L \times \frac{n}{2} + \frac{\phi kn^2}{4t^3} \times S - V - \frac{\phi n}{4t^3} \times kS + kV - 2T$$

$$C = M \times \frac{n}{3} + \frac{\phi kn^2}{9t^3} \times T - W - \frac{\phi n}{6t^3} \times kT + kW - 2V$$

$$D = N \times \frac{n}{4} + \frac{\phi kn^2}{16t^3} \times V - X - \frac{\phi n}{8t^3} \times kV + kX - 2W$$

&c.

$$\text{prodibit } v = \frac{\phi}{t^3} \times R - \frac{kS}{2} - \frac{2\phi kbn}{t^3} - Q \times s$$

$$+ A \times \sin. \frac{1}{n}s + B \times \sin. \frac{2}{n}s + C \times \sin. \frac{3}{n}s + D$$

$$\times \sin. \frac{4}{n}s +, \&c. + Z, \text{ et factâ hypothefi quòd fit}$$

$v = 0$ ubi angulus $PSQ = 0$, vel $r \times 180^\circ$, exprimente r unum ex numeris naturalibus 1, 2, 3, 4,

$$\&c. \text{ erit } Z = -\frac{\phi}{t^3} \times R - \frac{kS}{2} - \frac{2\phi kbn}{t^3} - Q \times s,$$

ac proinde $v = A \times \sin. \frac{1}{n}s + B \times \sin. \frac{2}{n}s + C \times \sin. \frac{3}{n}s + D \times \frac{4}{n}s +, \&c.$

Tùm, quia vis centripeta hîc excedere supponitur vim centrifugam, cùm contrarium suppositum fuerit in propositione secundâ, habetur — $\dot{x} = v \times \frac{xs}{\frac{1}{x} + U}$

five — $\dot{x} = vs$ proximè, et — $\frac{\dot{x}}{s} = v = K \times \frac{1}{n} \sin. \frac{1}{n}s + L \times \frac{2}{n} \sin. \frac{2}{n}s + M \times \frac{3}{n} \sin. \frac{3}{n}s + N \times \frac{4}{n} \sin. \frac{4}{n}s +, \&c.$

Unde factâ collatione terminorum hujus valoris velocitatis v cum terminis homologis valoris supra inventi, emergent

$$\begin{aligned} K &= -\frac{\phi}{t^3} \times \frac{n^2}{n^2-1} \times 2kR - \frac{2t^3}{k^2} \times n - \frac{1}{2} - kT \times n + \frac{1}{2} + S \\ L &= -\frac{\phi}{2t^3} \times \frac{n^2}{n^2-4} \times kS \times n - 1 - kV \times n + 1 + 2T \\ M &= -\frac{\phi}{3t^3} \times \frac{n^2}{n^2-9} \times kT \times n - \frac{3}{2} - kW \times n + \frac{3}{2} + 3V \\ N &= -\frac{\phi}{4t^3} \times \frac{n^2}{n^2-16} \times kV \times n - 2 - kX \times n + 2 + 4W \\ &\&c. \end{aligned}$$

atque ità patet hujusmodi quantitatum progressio. Innotescet igitur x , seu distantia planetæ P à Sole in quovis ejus cum planetâ Q aspectu.

Ut obtineatur planetæ P motus verus s , designet w motum medium, et cùm fit $\dot{w} = \frac{xs}{\frac{1}{x} + U}$,

tuantur

uantur valores quantitatum x , U , et sumptâ fluente, positis

$$F = 2nK + \frac{\phi k n^2}{t^3} \times \overline{R - \frac{t^3}{k^3} - \frac{T}{2}}$$

$$G = nL + \frac{\phi k n^2}{8t^3} \times \overline{S - V}$$

$$H = \frac{2nM}{3} + \frac{\phi k n^2}{18t^3} \times \overline{T - W}$$

$$I = \frac{nN}{2} + \frac{\phi k n^2}{32t^3} \times \overline{V - X}$$

&c.

proveniet $w = \overline{1 - 2Q - \frac{\phi k b n}{t^3} \times s} + F \times \text{fin. } \frac{1}{n}s + G \times \text{fin. } \frac{2}{n}s + H \times \text{fin. } \frac{3}{n}s + I \times \text{fin. } \frac{4}{n}s +$
&c. $+ Z$.

Et factâ hypothesi quod motus verus coincidat cum medio ubi est $\frac{1}{n}s$, seu angulus $PSQ = 0$, vel $= r \times 180^\circ$, exhibente r quemvis ex numeris 1, 2, 3, 4, &c. erit $Z = 2Q + \frac{\phi k b n}{t^3} \times s$; ac proinde, scriptis $\frac{1}{n}w$, $\frac{2}{n}w$, &c. pro $\frac{1}{n}s$, $\frac{2}{n}s$, &c. quia parùm admodùm differt motus verus à medio, habetur motus verus, five $s = w - F \times \text{fin. } \frac{1}{n}w - G \times \text{fin. } \frac{2}{n}w - H \times \text{fin. } \frac{3}{n}w - I \times \text{fin. } \frac{4}{n}w -$, &c. Q. E. I.

COROLL. I.

Designet jam planeta P Terram, Q Venerem, et quia posuimus esse distantiam mediocrem Terræ à Sole

Sole ad distantiam mediocrem Veneris à Sole ut 1 ad k , erit hîc $k = 0.72333$, atque $t = \sqrt{1 + kk} = 1.234182$. Item est $n = \frac{Q}{P - Q} = \frac{224.701}{365.2565 - 224.701} = 1.59866$. Quantitates b , R, S, T, &c. eisdem hîc retinent valores quos habebant in Coroll. I. Prop. II. Verùm, ut motuum Terrestrium accurata institueretur computatio, dignoscere necesse esset effectus aliquos ab actione Veneris provenientes, ex quibus derivare liceret vim attractivam istius planetæ, sed quia speciales hujusmodi effectus nulli, quantum noverimus, observationibus astronomicis explorati habentur, propterea vim Veneris nunc conjecturâ definiemus, ut inde inæqualitates in motu Telluris computatæ, atque cum observationibus astronomicis collatæ inservire posthac possint ad eandem vim certius determinandam. Itaque supponemus gravitatem in Solem esse ad gravitatem in Venerem, paribus distantiiis, ut 400000 ad 1, hoc est, esse $\phi = \frac{1}{400000}$. Qui tamen valor vis ϕ si major vel minor postea deprehensus fuerit, in eâdem ratione sequentes omnes determinaciones augendæ sunt, vel minuendæ, adeoque ad justam mensuram facillimè reducentur. Erunt igitur

$$\begin{array}{ll} K = - 0.00000575 & N = 0.00000090 \\ L = 0.00001643 & O = 0.00000039 \\ M = 0.00000259 & O' = 0.00000022, \text{ \&c.} \end{array}$$

Indeque colliguntur

$$\begin{array}{ll} F = - 0.00002459 & I = 0.00000105 \\ G = 0.00002795 & I' = 0.00000042 \\ H = 0.00000345 & \text{ \&c.} \end{array}$$

atque reductis quantitatibus F, G, H, &c. in partes circuli,

circuli, tandem habetur $s = w + 5''.07 \times \sin. \frac{1}{n} w$
 $- 5''.76 \times \sin. \frac{2}{n} w - 0''.71 \times \sin. \frac{3}{n} w - 0''.22$
 $\times \sin. \frac{4}{n} w -$, &c. ubi s denotat motum Terræ verum,
 w motum medium, et $\frac{1}{n} w$ angulum PSQ, five dif-
ferentiam longitudinum heliocentricarum Terræ et
Veneris.

Inde computatur sequens tabula exhibens æqua-
tionem motûs Solis pro variâ distantîâ Veneris à Terrâ
quam metitur angulus PSQ, five pro variâ differentiâ
longitudinum heliocentricarum Terræ et Veneris quam
metitur arcus circuli maximi inter Terram et Venerem
interjectus et secundum seriem signorum à loco Terræ
computatus.

Diff. long. hel. Terræ et Ven.	Æquatio motûs Solis.	Diff. long. hel. Terræ et Ven.	Æquatio motûs Solis.
°	°	°	°
Sig. o. 0	— 0	Sig. VI. 0	— 0
10	1.6	10	2.6
20	2.8	20	5.0
30	3.4	30	7.0
Sig. I. 10	3.1	Sig. VII. 10	8.4
20	2.1	20	9.1
30	0.4	30	9.2
Sig. II. 10	+ 1.6	Sig. VIII. 10	8.6
20	3.8	20	7.5
30	5.8	30	5.8
Sig. III. 10	7.5	Sig. IX. 10	3.8
20	8.6	20	1.6
30	9.2	30	+ 0.4
Sig. IV. 10	9.1	Sig. X. 10	2.1
20	8.4	20	3.1
30	7.0	30	3.4
Sig. V. 10	5.0	Sig. XI. 10	2.8
20	2.6	20	0.6
30	0.	30	0

COROLL. II.

Si tellus gravitate suâ in Solem in circulo revolvi posse supponatur, adveniente Veneris actione variari debere distantiam ejus à Sole patet ex hac propositione. Esto angulus $\frac{1}{n}s$, seu $PSQ = 90^\circ$, vel 270° , atque æquatio generalis $x = 1 - Q + K \cos. \frac{1}{n}s + L \cos. \frac{2}{n}s + M \cos. \frac{3}{n}s +$, &c. in hanc abit $x = 0.9999693$; et si fit $PSQ = 180^\circ$, fit $x = 1.0000053$. Unde si distantia Terræ à Sole, ubi versatur in conjunctione cum Venere, } 10000000
ponatur - - - - - }
In quadraturis cum Venere erit ipsius distantia - - - - - } 9999693
Atque in oppositione - - - - - } 10000053

PROPOSITIO VII. PROBLEMA.

In systemate duorum planetarum in circulis circa Solem revolventium, motum nodorum orbis planetæ exterioris in plano orbis planetæ interioris investigare.

Esto P locus planetæ exterioris (Fig. 5.) in orbe suo PN, SQ recta conjungens Solem et planetam interiorem, et dicatur c finis inclinationis duorum orbium ad se invicem ad radium 1, atque per propositionem quintam est $\frac{\phi k}{x^3} - \frac{\phi}{k^2}$ vis qua planeta P amovetur ab orbe suo secundum directionem parallelam rectæ SQ, hujusque vis ea pars quæ perpendiculariter

agit in planum orbis PN, per simile ratiocinium quo
uti sumus in Prop. III. prodit æqualis $c \times \sin. QN$

$\times \frac{\phi k}{z^3} - \frac{\phi}{k^2}$, et motus interfectionis plani orbis PN

cum plano orbis QN fit $\frac{\phi k}{z^3} - \frac{\phi}{k^2} \times \sin. PN \times \sin. QN$

$\times Pp$ quo tempore planeta P describit in orbe suo
arcum quàm minimum Pp.

Deinde si designaverit D locum planetæ P ubi ver-
satur in conjunctione cum planetâ interiore, et ponan-
tur $DP = s$, $Pp = i$, $DN = a$, erit $PN = s + a$,

$QN = s + \frac{1}{n}s + a$ quamproximè, atque $\sin. PN$

$\times \sin. QN = \frac{1}{2} \cos. \frac{1}{n}s - \frac{1}{2} \cos. 2s + \frac{1}{n}s + 2a$.

Unde, calculum profequendo uti in propositione
tertiâ, motus nodorum factus, quo tempore planeta
P à loco conjunctionis D discedens descripserit in
orbe suo arcum quemlibet DP, exprimetur per

$\frac{\phi kn}{2t^3} \sin \frac{S}{2n}s + R - \frac{t^3}{k^3} + \frac{T}{2} \times \sin. \frac{1}{n}s + \frac{S+V}{4} \sin. \frac{2}{n}s$

$+ \frac{T+W}{6} \sin. \frac{3}{n}s + \frac{V+X}{8} \sin. \frac{4}{n}s +$, &c.

$+ \frac{\phi kn}{2t^3} \sin Z \times \sin. 2a - R - \frac{t^3}{k^3} \times \frac{1}{2n+1} \sin. 2s + \frac{1}{n}s + 2a$

$- \frac{S}{2} \times \frac{1}{2n} \sin. 2s + 2a - \frac{S}{2} \times \frac{1}{2n+2} \sin. 2s + \frac{2}{n}s + 2a$

$- \frac{T}{2} \times \frac{1}{2n-1} \sin. 2s - \frac{1}{n}s + 2a - \frac{T}{2} \times \frac{1}{2n+3} \sin.$

$2s + \frac{3}{n}s + 2a - \frac{V}{2} \times \frac{1}{2n-2} \sin. 2s - \frac{2}{n}s + 2a$

$$\begin{aligned}
 & -\frac{V}{2} \times \frac{1}{2n+4} \sin. 2s + \frac{4}{n}s + 2a - \frac{W}{2} \times \frac{1}{2n-3} \sin. \\
 & 2s - \frac{3}{n}s + 2a - \frac{W}{2} \times \frac{1}{2n+5} \sin. 2s + \frac{5}{n}s + 2a, \&c. \\
 & \text{existente } Z = 2n + 1 \text{ in } R - \frac{t^3}{k} \times \frac{1}{2n+1!^2} + \frac{S}{2n \times 2n+2} \\
 & + \frac{T}{2n-1 \times 2n+3} + \frac{V}{2n-2 \times 2n+4} + \frac{W}{2n-3 \times 2n+5} \\
 & +, \&c. \text{ In quibus seriebus manifesta est terminorum} \\
 & \text{progressio. Q. E. I.}
 \end{aligned}$$

COROLL.

Hinc in conjunctionibus expressio motûs nodi evadit

$$\begin{aligned}
 & \frac{\phi k}{2t^3} \times \frac{S}{2}s - nZ \times \sin. 2s + 2a - \sin. 2a. \text{ Hic} \\
 & \text{que est motus nodi factus quo tempore planetæ P} \\
 & \text{et Q à conjunctione procedentes ad conjunctionem} \\
 & \text{quamvis aliam pervenerint, exhibente } s \text{ arcum à pla-} \\
 & \text{netâ P in suâ orbitâ intereâ descriptum. Terminus} \\
 & \frac{\phi k}{2t^3} \times \frac{S}{2}s \text{ exprimit motum nodi medium, et terminus} \\
 & \text{alter } \frac{\phi kn}{2t^3} Z \times \sin. 2s + 2a - \sin. 2a \text{ indicat æqua-} \\
 & \text{tionem } \textit{periodicam generalem}; \text{ vel etiam, si conjunctio} \\
 & \text{illa à qua desumitur computationis initium, fieri sup-} \\
 & \text{ponatur in nodo, vel propè ad nodum, æquatio pe-} \\
 & \text{riodica generalis fit } \frac{\phi kn}{2t^3} Z \times \sin 2s.
 \end{aligned}$$

Designet jam planeta P Terram, Q Venerem, eritque post unam revolutionem synodicam, id est, post revolutionem Veneris ad Terram, $\frac{1}{n}s = 360^\circ$,

U u 2

proindeque

proindeque $s = n \times 360^\circ = 575^\circ 31'$. Quare motus nodi medius huic temporis spatio congruens fit $\frac{\phi k n}{4t^3} S \times 360^\circ$, qui imminutus in ratione revolutionis Terræ circa Solem ad ejusdem revolutionem ad Venerem, hoc est, in ratione 1 ad n , evadit $\frac{\phi k}{4t^3} S \times 360^\circ = 5''.20$, motus scilicet nodi medius annuus quo regreditur intersectio planorum orbium Terræ ac Veneris; atque hic motus spatio centum annorum fit $8' 40''$.

In computo æquationis *periodicæ generalis* $\frac{\phi k n}{2t^3} Z \times \sin. 2s$, advertendum est omnes terminos, ex quibus componitur valor quantitatis Z , eisdem hîc esse ac in Prop. III. præter terminum primum $R - \frac{t^3}{k^3} \left(\times \frac{1}{2n+1} \right)$ qui ob diversum valorem quantitatum t et k diversus est. Hîc igitur provenit $Z = 31.59$, adeoque $\frac{\phi k n}{2t^3} Z \times \sin. 2s = 5'' \times \sin. 2s$; unde patet æquationem hanc nunquam superare $5''$. Motus igitur nodi verus, nimirum $\frac{\phi k}{2t^3} \times \frac{S}{2} s - nZ \times \sin. 2s$, peractâ unâ revolutione synodicâ post conjunctionem factam in nodo, evadit $8'.3 - 5'' \times \sin. 71^\circ.2'$, quia tunc est $\sin. 2s = \sin. 2 \times 575^\circ.31' = \sin. 71^\circ.2'$; et per ratiocinium simile ei, quod in Coroll. II. Prop. III. usurpatum est, constabit $8''.3 - 5''.8 \times \cos. 2r - 1 \times 35^\circ.31'$ exprimere regressum nodi factum tempore illius revolutionis synodicæ, cujus lo-

cum in ferie revolutionum indicat numerus r . Hinc computatur tabula sequens quæ exhibet regressum nodi orbitæ Terrestris in plano orbis Veneris pro duodecim sigillatim revolutionibus synodicis quæ proximè sequuntur conjunctionem Terræ et Veneris factam in nodo, vel proximè ad nodum.

In revol. synod.	Regressus nodi Ter.	In revol. synod.	Regressus nodi Ter.
	//		//
1	4	7	9
2	10	8	14
3	14	9	11
4	10	10	4
5	4	11	3
6	3	12	9

Patet autem æquationem *periodicam specialem*, nempe $5''.8 \times \cos. 2r - 1 \times 35^\circ. 31'$, ubi maxima est, evadere $5''.8$, et regressum nodi in quavis revolutione Terræ ad Venerem non assurgere ultra $14''$, nec minui citra $2''\frac{1}{2}$.

PROPOSITIO VIII. PROBLEMA.

Iisdem positis, variationem inclinationis orbis planetæ exterioris ad planum orbis planetæ interioris determinare.

Designet I variationem inclinationis factam quo tempore planeta P describit arcum quàm minimum Pp ,

Pp, et N motum nodi eodem tempore confectum, ac per ratiocinium omnino simile ei quod adhibitum est in propositione quartâ habetur $I = N \times \frac{c \times \cos. PN}{\sin. PN}$:

sed per propositionem præcedentem est $N = \frac{\phi k}{z^3} - \frac{\phi}{k^2}$
 $\times \sin. PN \times \sin. QN \times Pp$, adeoque fit $I = \frac{\phi k}{z^3} - \frac{\phi}{k^2}$
 $\times c \times \cos. PN \times \sin. QN \times Pp$.

Unde, cùm hîc fit $PN = s + a$, $QN = s + \frac{1}{n}s + a$, proindeque $\cos. PN \times \sin. QN = \frac{1}{2} \sin. \frac{1}{n}s + \frac{1}{2} \sin. 2s + \frac{1}{n}s + 2a$, sumptâ fluente prodit variatio inclinationis genita, quo tempore planeta descripserit in orbe suo arcum quemlibet DP à loco conjunctionis D, æqualis $\frac{\phi ckn}{2t^3}$ in $R - \frac{t^3}{k^3} - \frac{T}{2}$
 $\times \sin. \text{verf. } \frac{1}{n}s + \frac{S-V}{4} \sin. \text{verf. } \frac{2}{n}s + \frac{T-W}{6} \sin. \text{verf. } \frac{3}{n}s + \frac{V-X}{8} \sin. \text{verf. } \frac{4}{n}s +, \&c. + \frac{\phi ckn}{2t^3}$ in
 $- Z \times \sin. \text{verf. } 2a + R - \frac{t^3}{k^3} \times \frac{1}{2n+1} \sin. \text{verf. } 2s + \frac{1}{n}s + 2a + \frac{S}{2} \times \frac{1}{2n} \sin. \text{verf. } 2s + 2a + \frac{S}{2}$
 $[\times \frac{1}{2n+2} \sin. \text{verf. } 2s + \frac{2}{n}s + 2a + \frac{T}{2} \times \frac{1}{2n-1} \sin. \text{verf. } 2s - \frac{1}{n}s + 2a + \frac{T}{2} \times \frac{1}{2n+3} \sin. \text{verf. } 2s + \frac{3}{n}s + 2a + \frac{V}{2} \times \frac{1}{2n-2} \sin. \text{verf. } 2s - \frac{2}{n}s + 2a$
 $+ +$

$\div \frac{V}{2} \times \frac{1}{2n+4}$ fin. verf. $2s \div \frac{4}{n}s \div 2a$, &c. Eundem hîc habet valorem quantitas Z ac in propositione præcedente. Q. E. I.

C O R O L L.

Ubi angulus PSQ est nullus, vel multiplex anguli 360° , id est, ubi planetæ versantur in conjunctione, variatio inclinationis genita generatim est $\frac{\phi c k n}{2t^3} Z$
 \times fin. verf. $2s \div 2a -$ fin. verf. $2a$ quæ, si ponatur arcus $DN = a = 0$, fit $\frac{\phi c k n}{2t^3} Z \times$ fin. verf. $2s$.

Atque hoc est decrementum inclinationis orbis planetæ P ad orbem planetæ Q factum in qualibet serie revolutionum ad conjunctionem, initio sumpto à conjunctione factâ in nodo, vel prope ad nodum, et designante s arcum intereà à planetâ P in orbe suo descriptum.

Si inde computetur decrementum inclinationis orbis Terrestris supra planum orbitæ Veneris factum post quocumque revolutiones Veneris ad Terram, fiet $\frac{\phi c k n}{2t^3} Z \times$ fin. verf. $2s = 0''.3 \times$ fin. verf. $2s$, adeoque hoc decrementum, ubi maximum evadit, non superat $0''.6$, ac proinde in omni casu negligi potest.